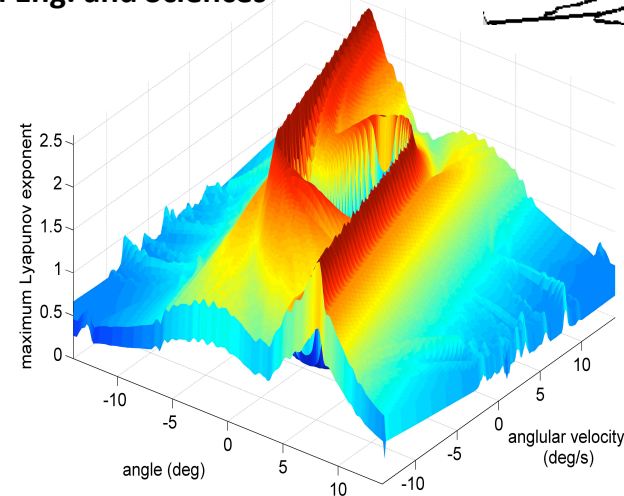
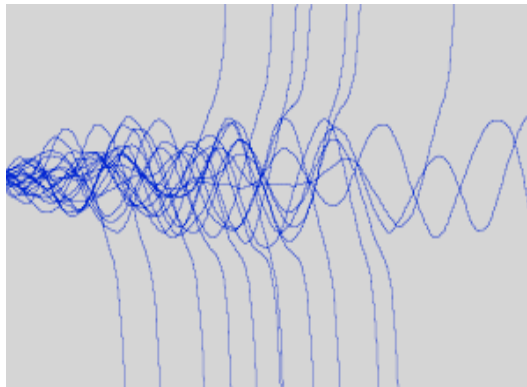
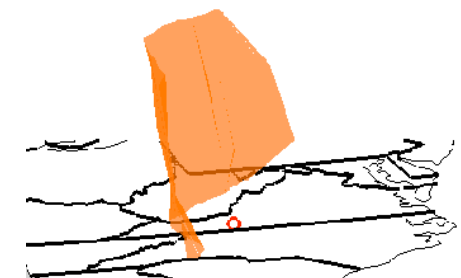
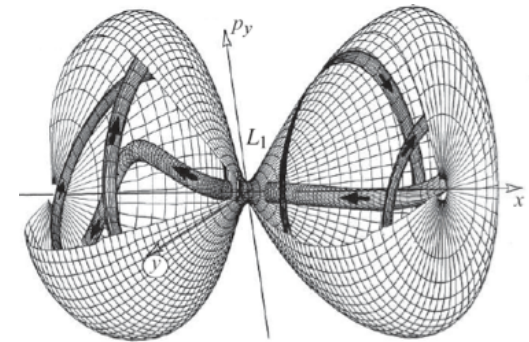


Dynamical boundaries in a variety of mechanical systems

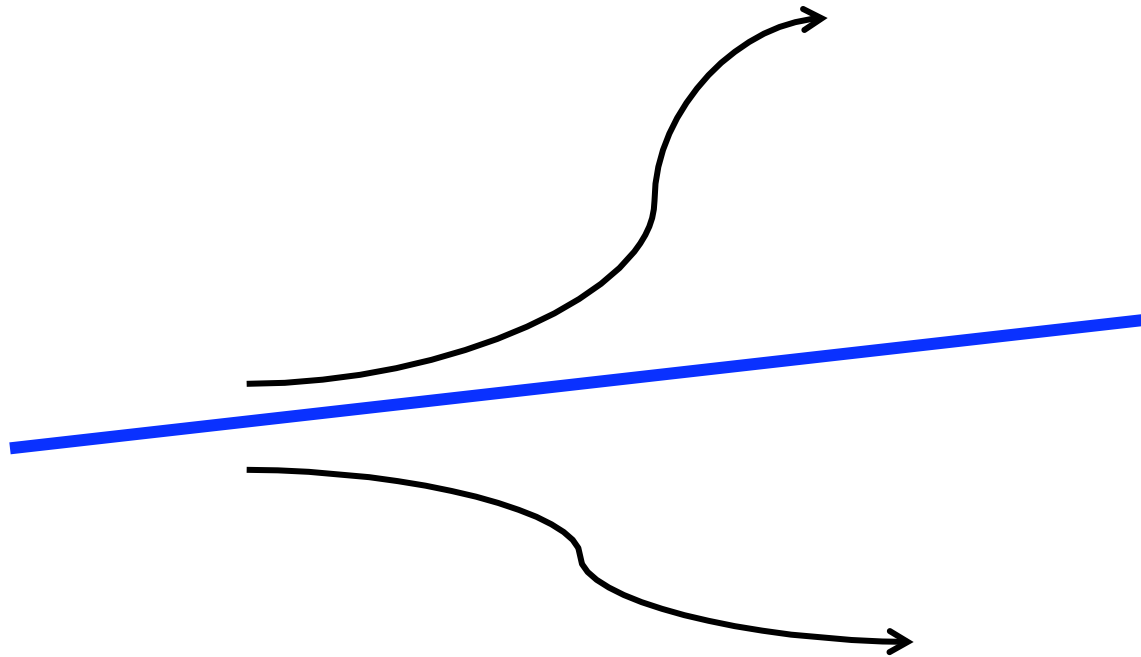
Shane D. Ross

Engineering Science and Mechanics, Virginia Tech

Virginia Tech-Wake Forest Univ. School of Biomedical Eng. and Sciences

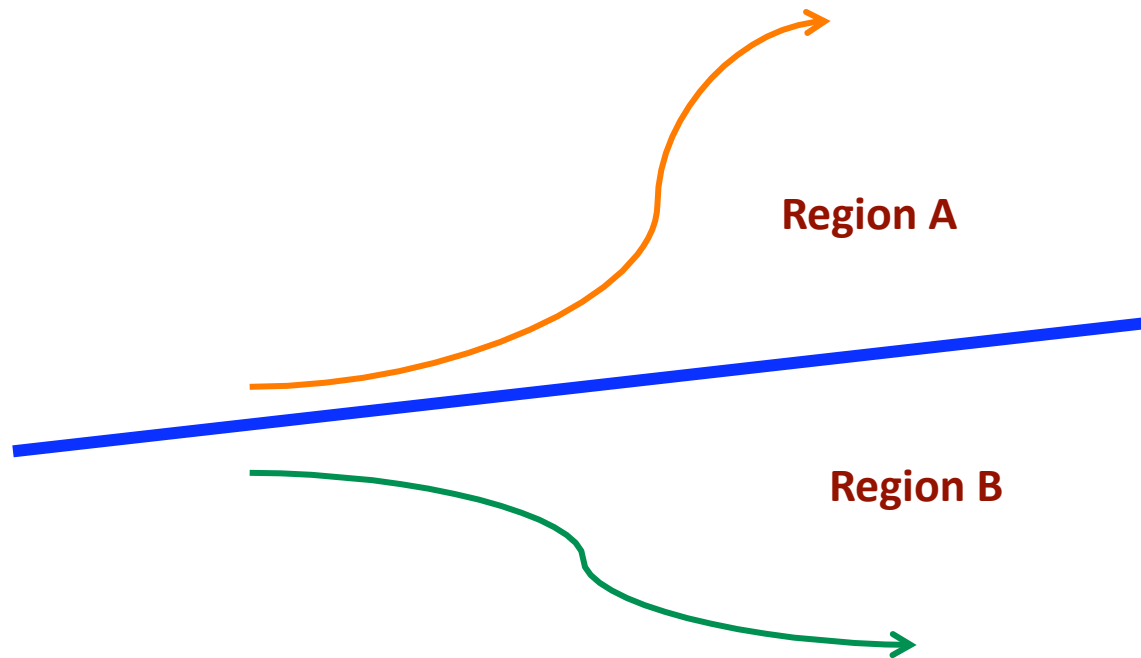


Separatrices: dynamical boundaries



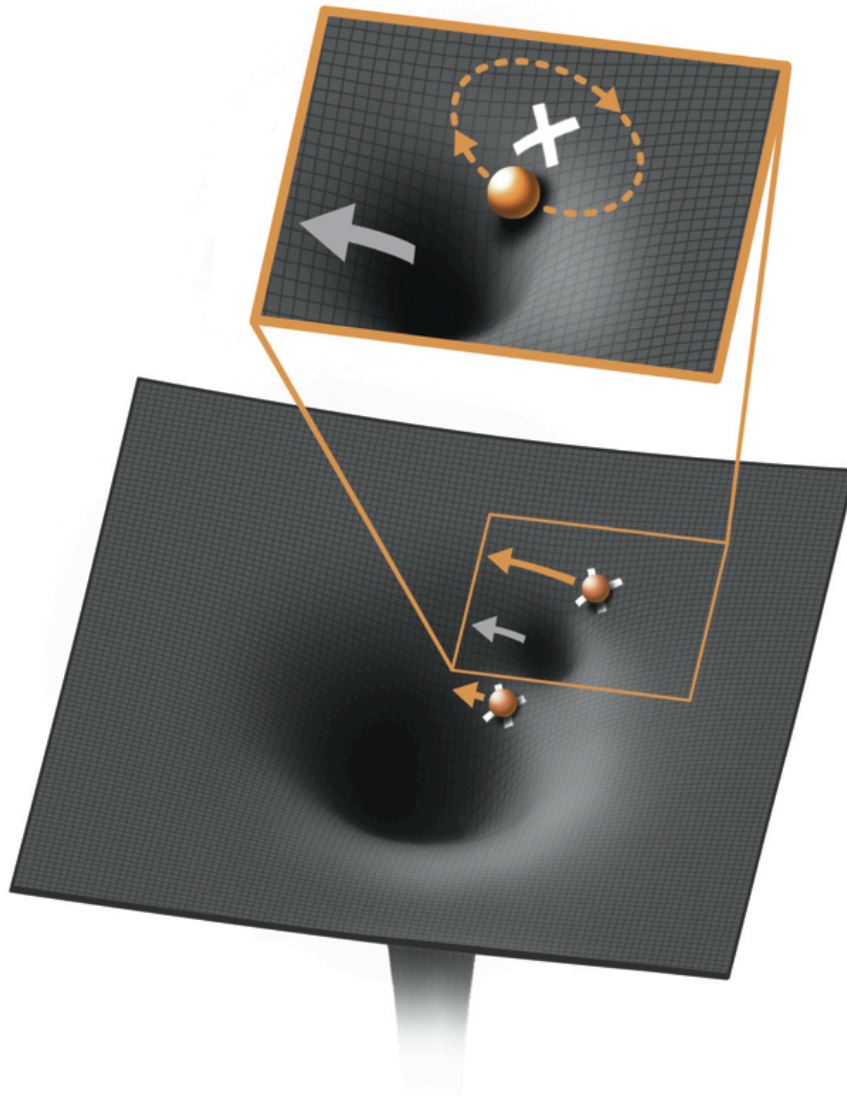
**Transport barriers in state space
separating qualitatively different kinds of behavior**

Separatrices: dynamical boundaries



**Transport barriers in state space
separating qualitatively different kinds of behavior**

Separatrices: high dimensions



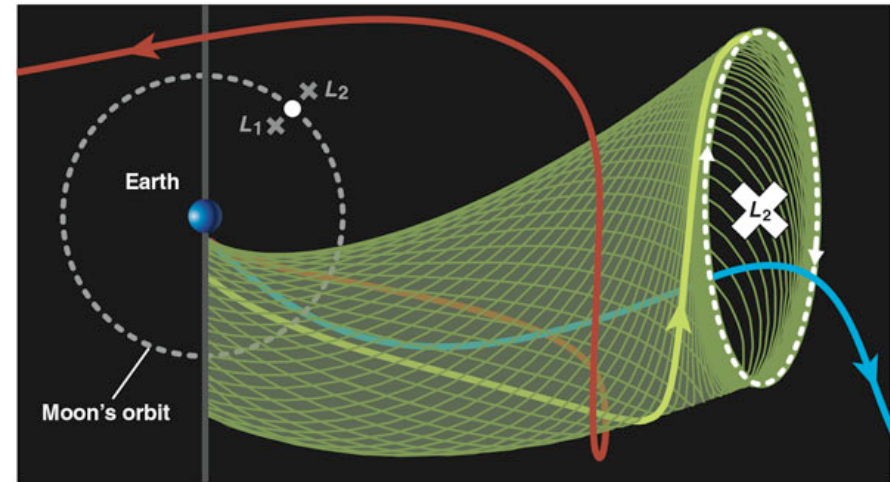
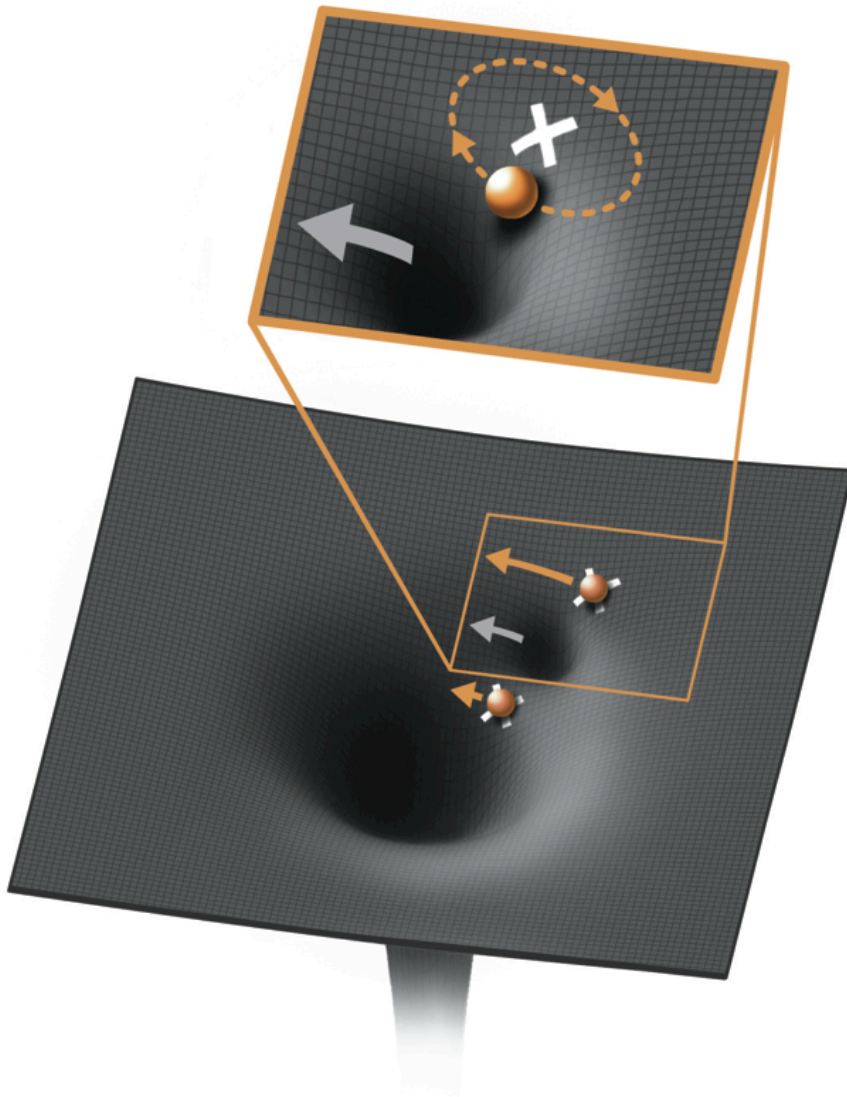
System with many basins, not necessarily attracting sets or attractors

Potential surface with several minima (“bowls”) separated by ridges

Basins are “almost-invariant structures”

System known analytically: vector field or map

Separatrices: high dimensions



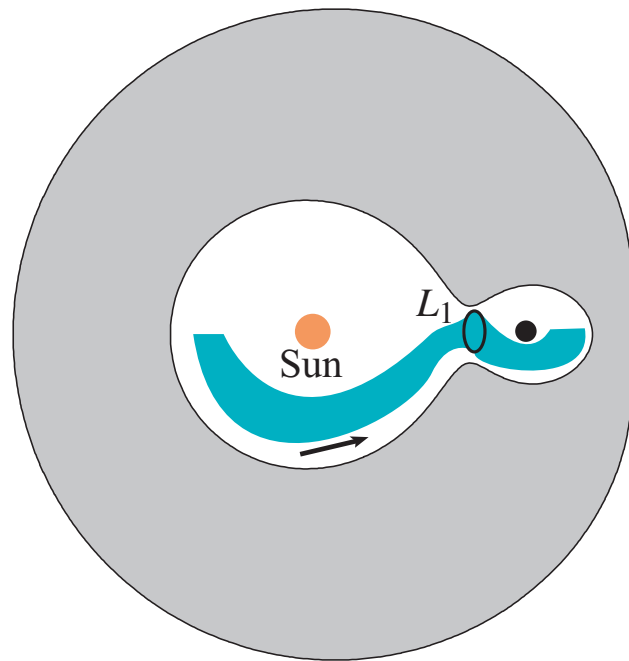
Small body in solar system: Transport from one basin to another controlled by high dimensional separatrix surfaces

Geometrically tubes in this case

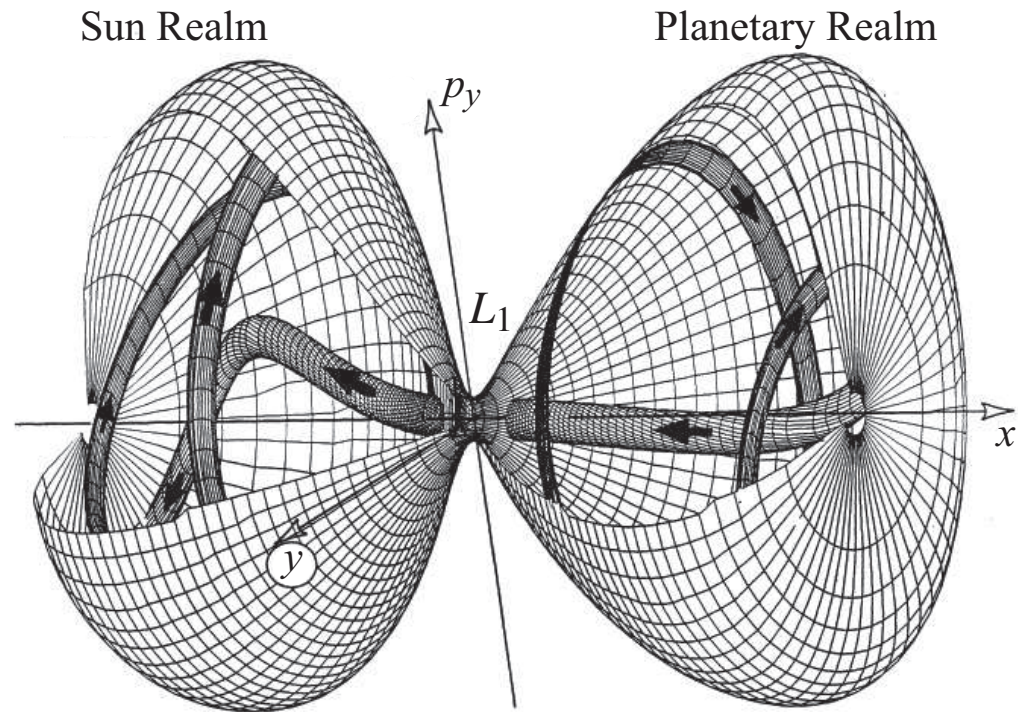
Tubes are attached to practically unobservable periodic orbits or other bound orbits

Realms and tubes

- Planetary and sun realms connected by tubes¹



Position Space

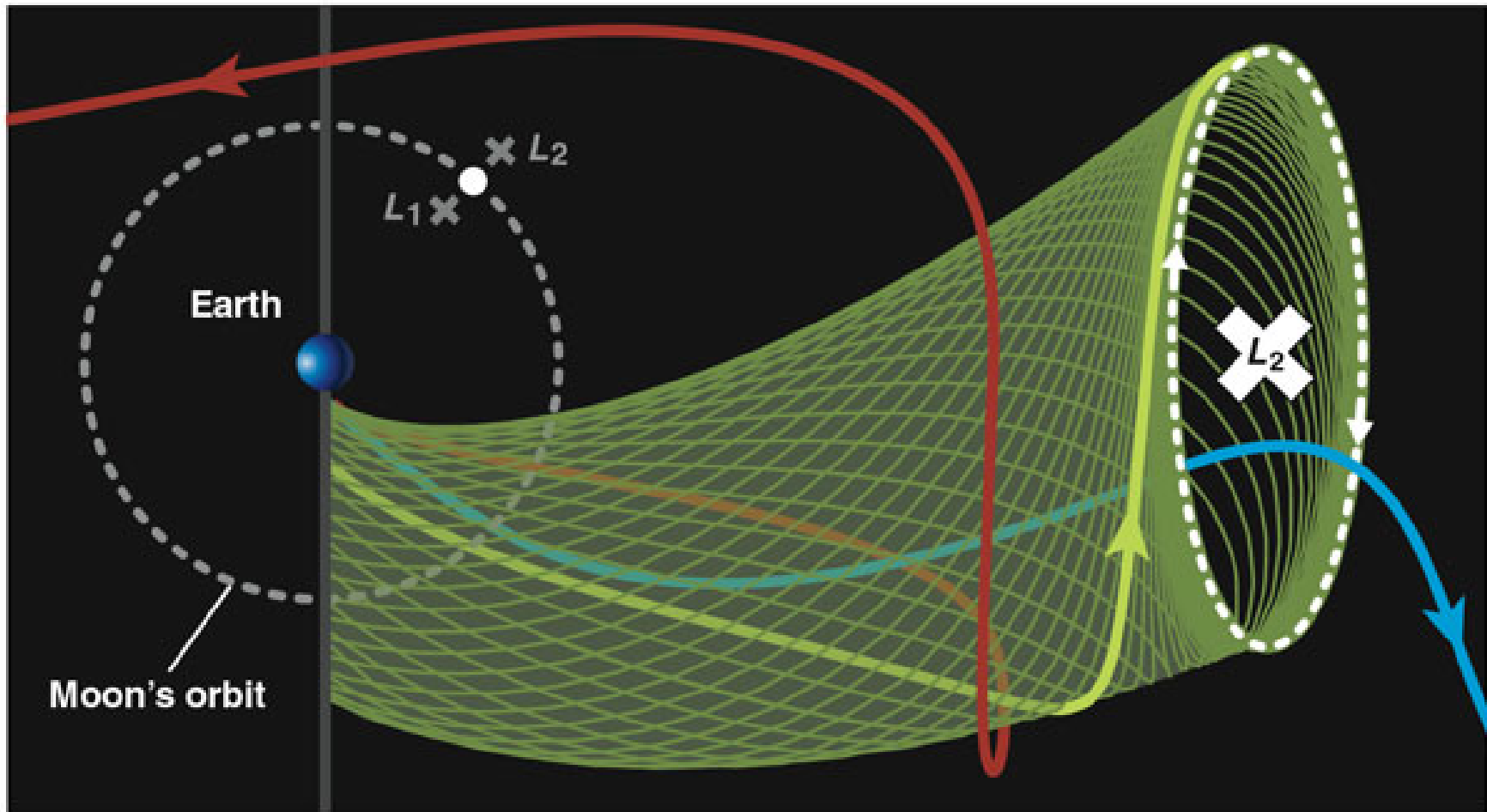


Phase Space (Position + Velocity)

¹Conley & McGehee, 1960s, found these locally, speculated use for “low energy transfers”

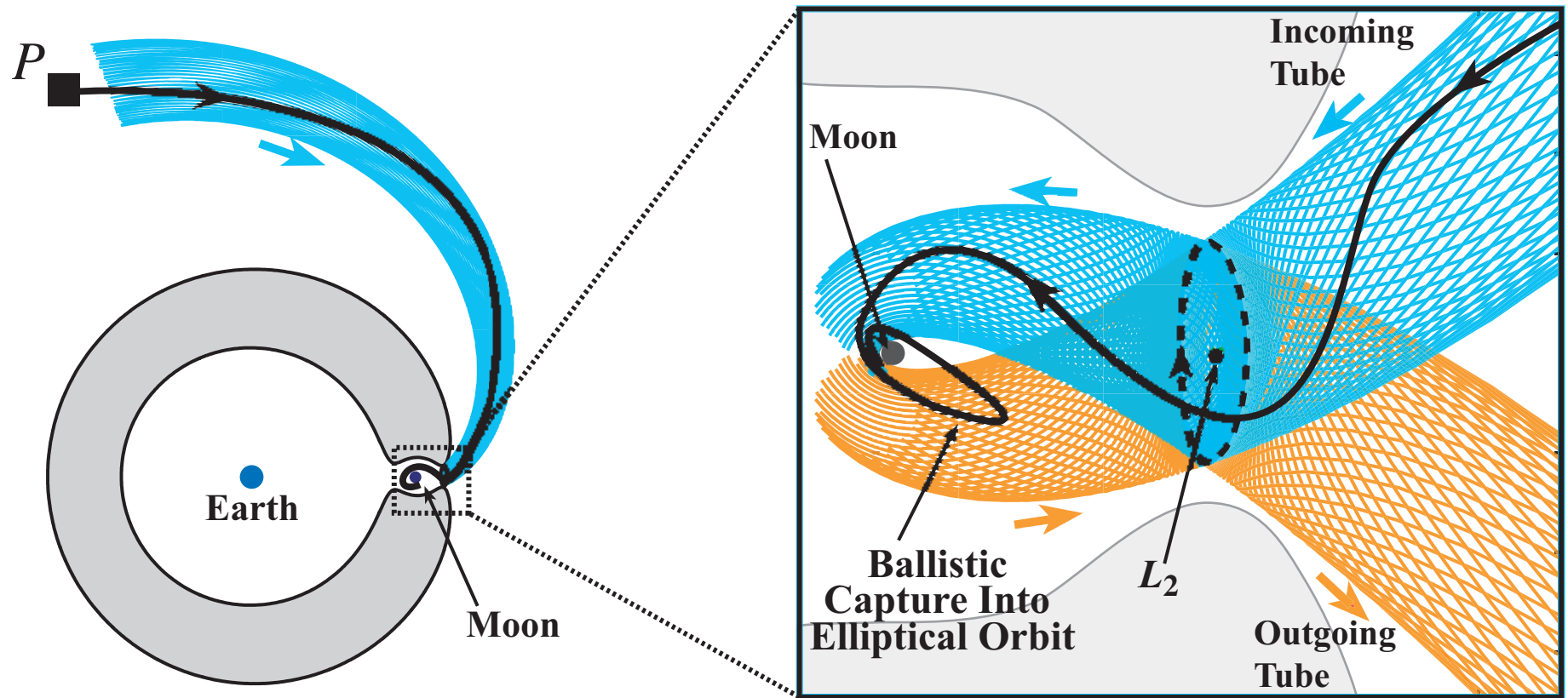
Transport between realms

- Asymptotic orbits form **4D invariant manifold tubes** ($S^3 \times \mathbb{R}$), separatrices in 5D energy surface²



²Ross [2006] The interplanetary transport network, *American Scientist*

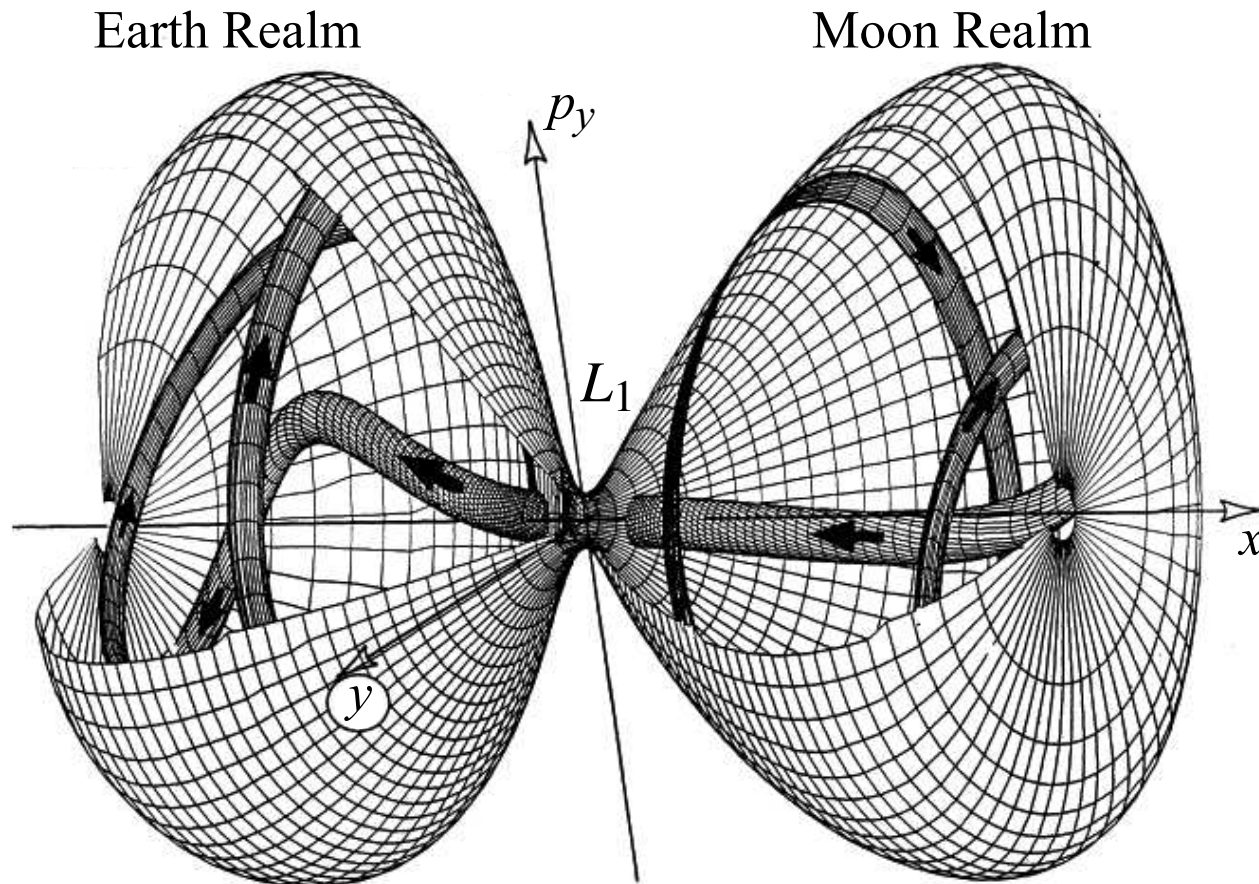
Transport between realms



□ Tubes in phase space

- Objects mediating transport through bottlenecks

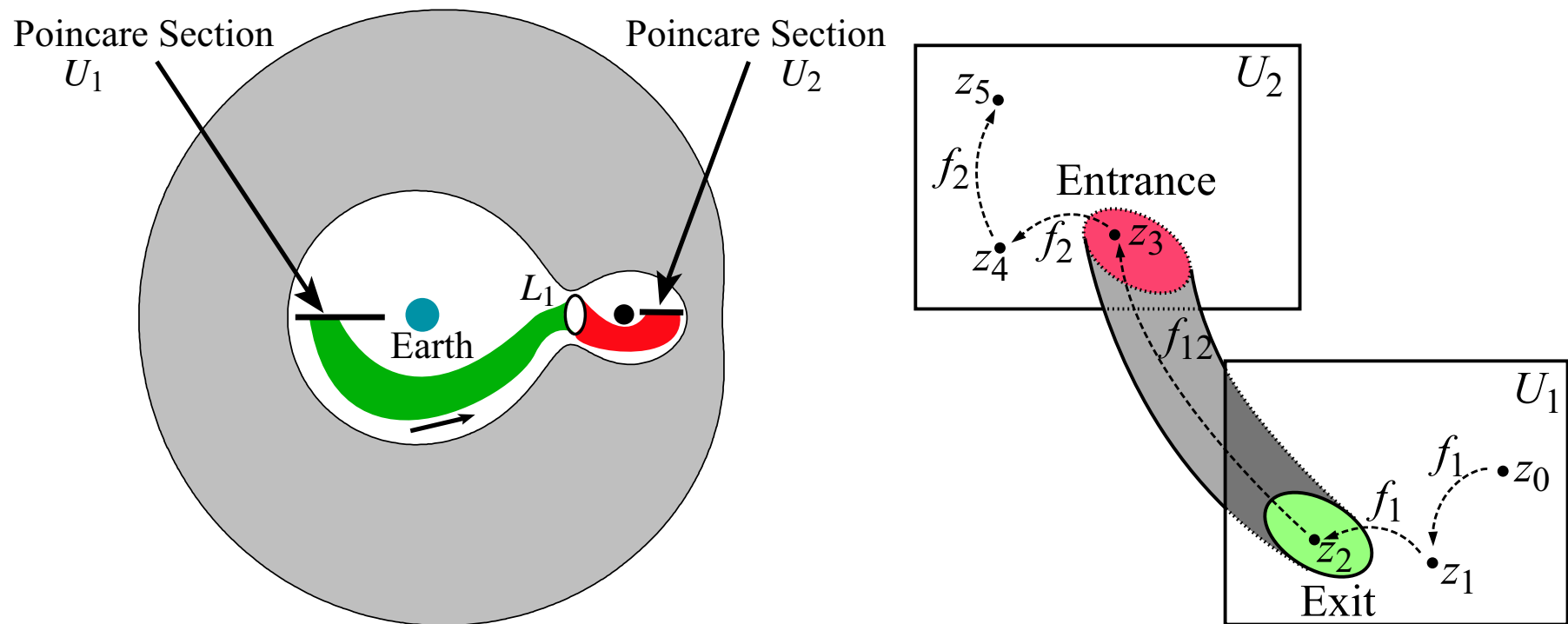
Tube dynamics



- **Tube dynamics:** All motion between realms connected by bottlenecks must occur through the interior of tubes

Multi-scale dynamics

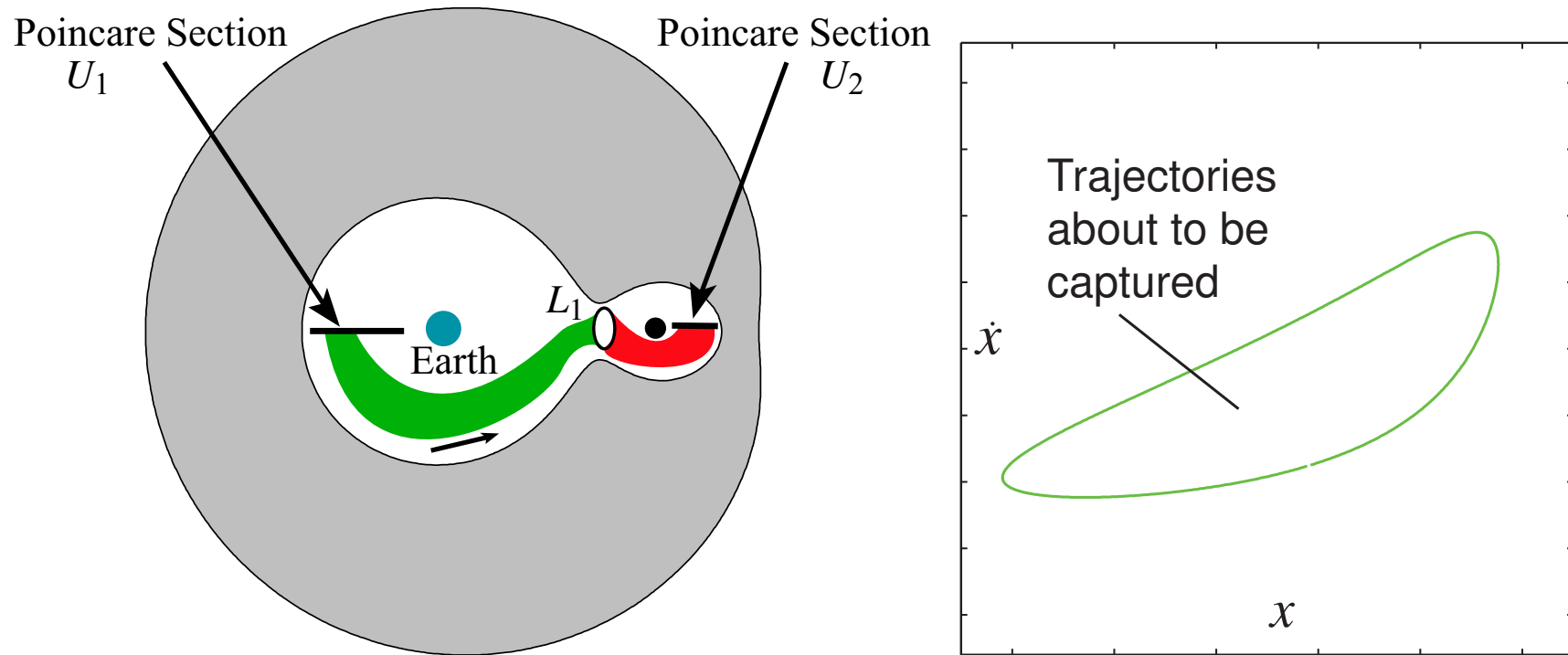
- Slices of energy surface: Poincaré sections U_i
- Tube dynamics: evolution **between** U_i
- What about evolution **on** U_i ?



Some remarks on tube dynamics

- Tubes are general; consequence of rank 1 saddle
 - saddle \times center $\times \dots \times$ center – e.g., ubiquitous in chemistry
- Tubes persist
 - in presence of additional massive body
 - when primary bodies' orbit is eccentric

Tubes in elliptic restricted 3-body problem



Consider first cut of stable manifold of L_1 NHIM

Tubes in elliptic restricted 3-body problem

Gawlik, Marsden, Du Toit, Campagnola [2008] “Lagrangian coherent structures in the planar elliptic restricted three-body problem,” submitted to *Celestial Mechanics and Dynamical Astronomy*.

Tubes in elliptic restricted 3-body problem

Gawlik, Marsden, Du Toit, Campagnola [2008] “Lagrangian coherent structures in the planar elliptic restricted three-body problem,” submitted to *Celestial Mechanics and Dynamical Astronomy*.

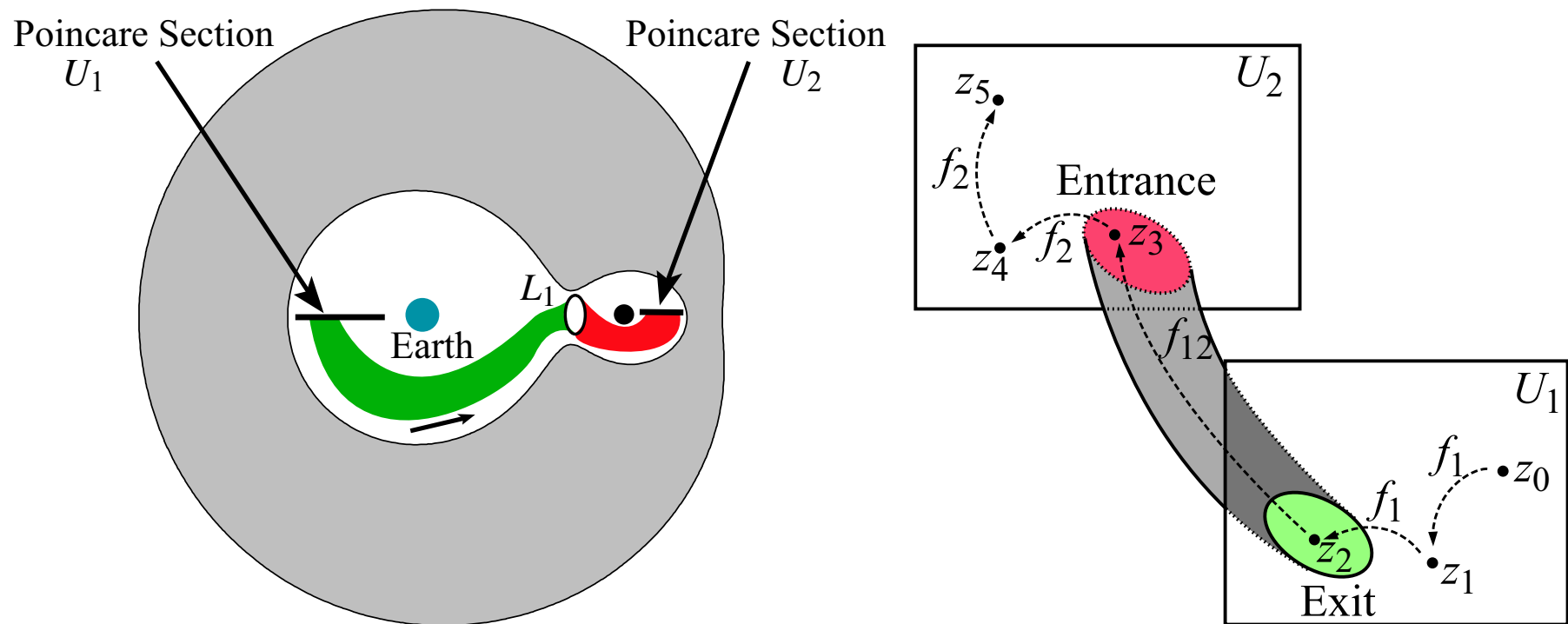
Some remarks on tube dynamics

- Tubes are general; consequence of rank 1 saddle
 - saddle \times center $\times \dots \times$ center
 - e.g., ubiquitous in chemistry
- Tubes persist
 - in presence of additional massive body
 - when primary bodies' orbit is eccentric
- Observed in the solar system (e.g., Oterma)
- Even on galactic and atomic scales!

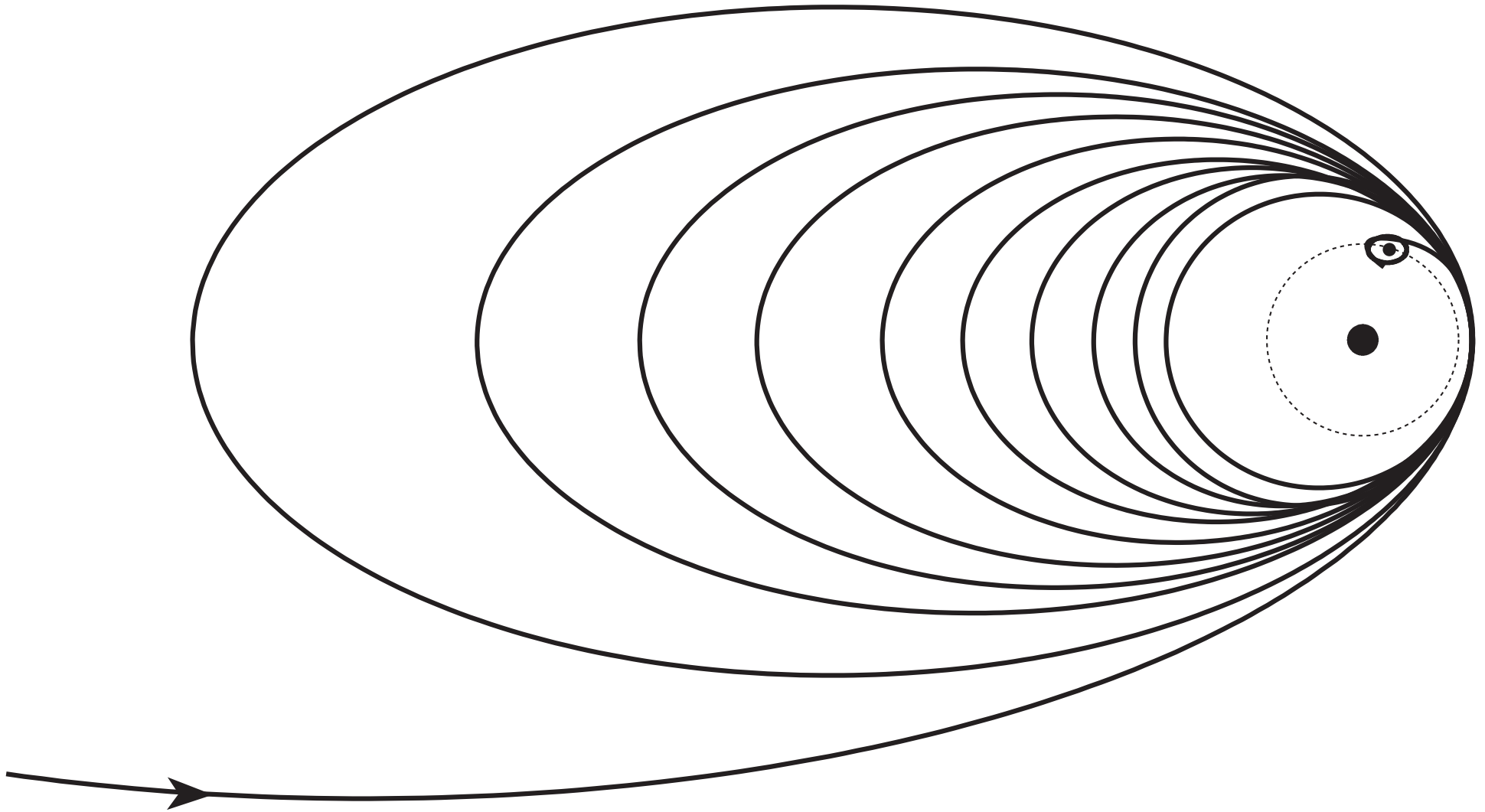
Koon, Lo, Marsden, & Ross [2000], Gómez, Koon, Lo, Marsden, Masdemont, & Ross [2004], Yamato & Spencer [2003], Wilczak & Zgliczyński [2005], Ross & Marsden [2006], Gawlik, Marsden, Du Toit, Campagnola [2008], Combes, Leon, Meylan [1999], Heggie [2000], Romero-Gómez, et al. [2006,2007,2008]

Multi-scale dynamics

- Slices of energy surface: Poincaré sections U_i
- Tube dynamics: evolution **between** U_i
- \longrightarrow What about evolution **on** U_i ? \longleftarrow



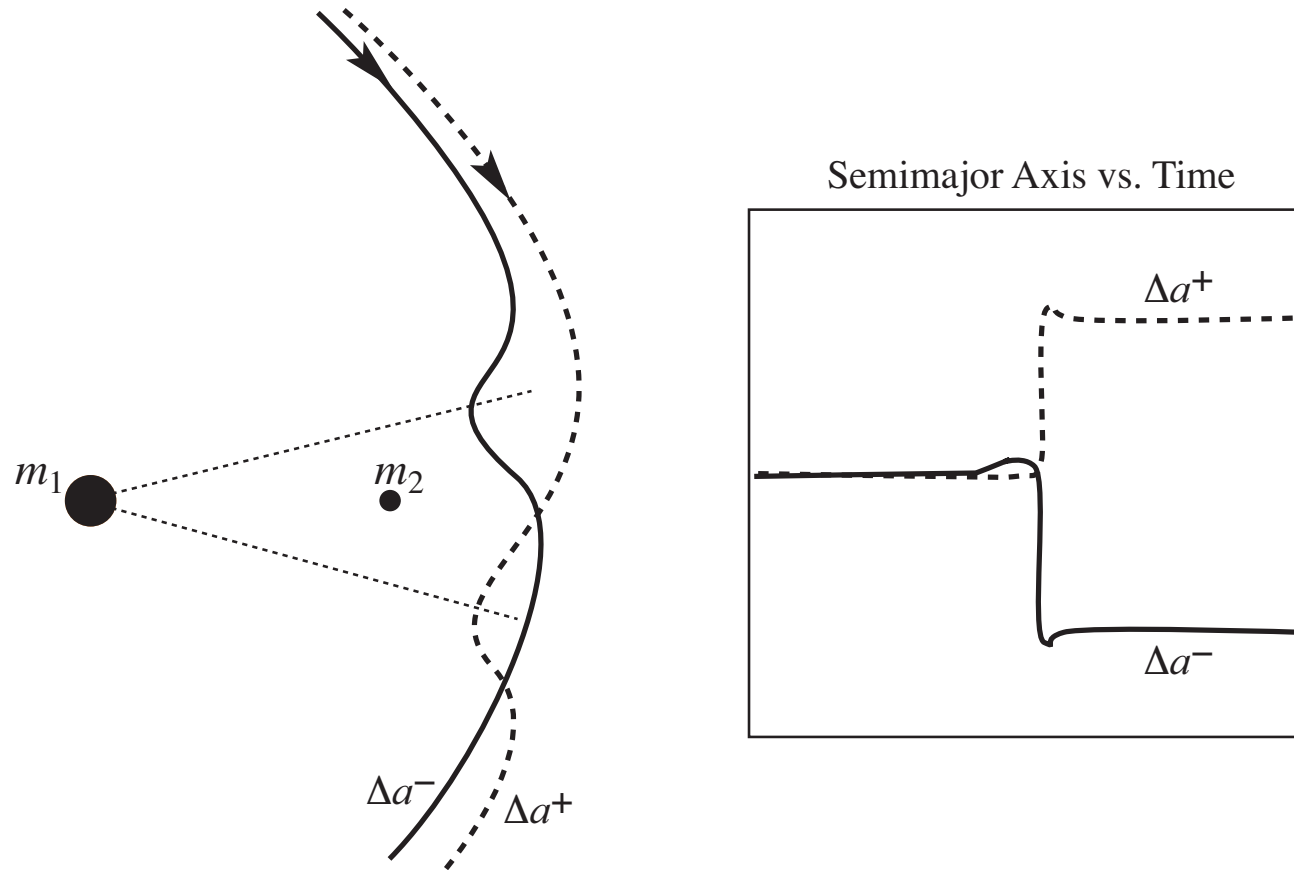
Infinity to capture about small companion in binary pair?



□ After **consecutive gravity assists**, large orbit changes

Kicks at periapsis

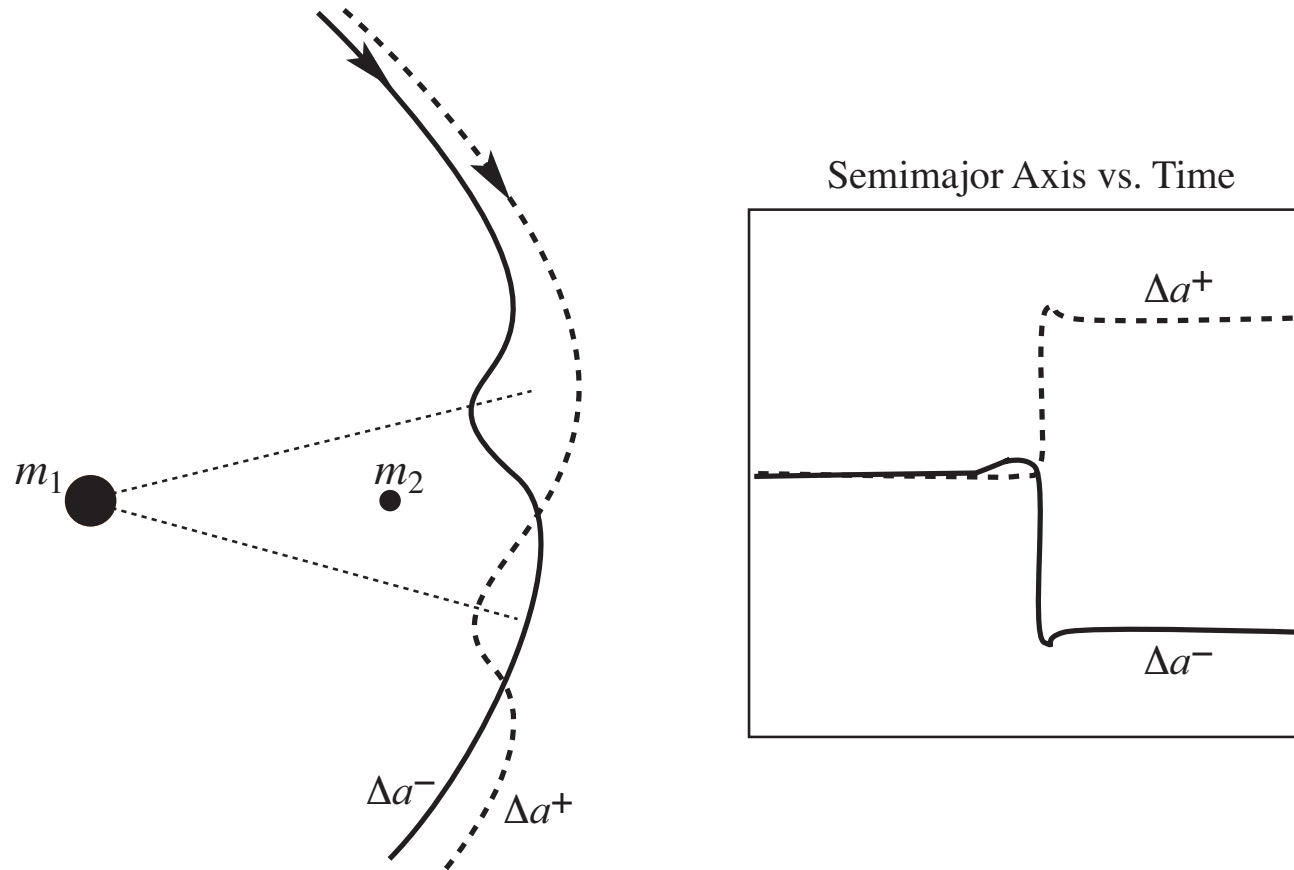
- Key idea: model particle motion as “kicks” at periapsis



In rotating frame where m_1, m_2 are fixed

Kicks at periapsis

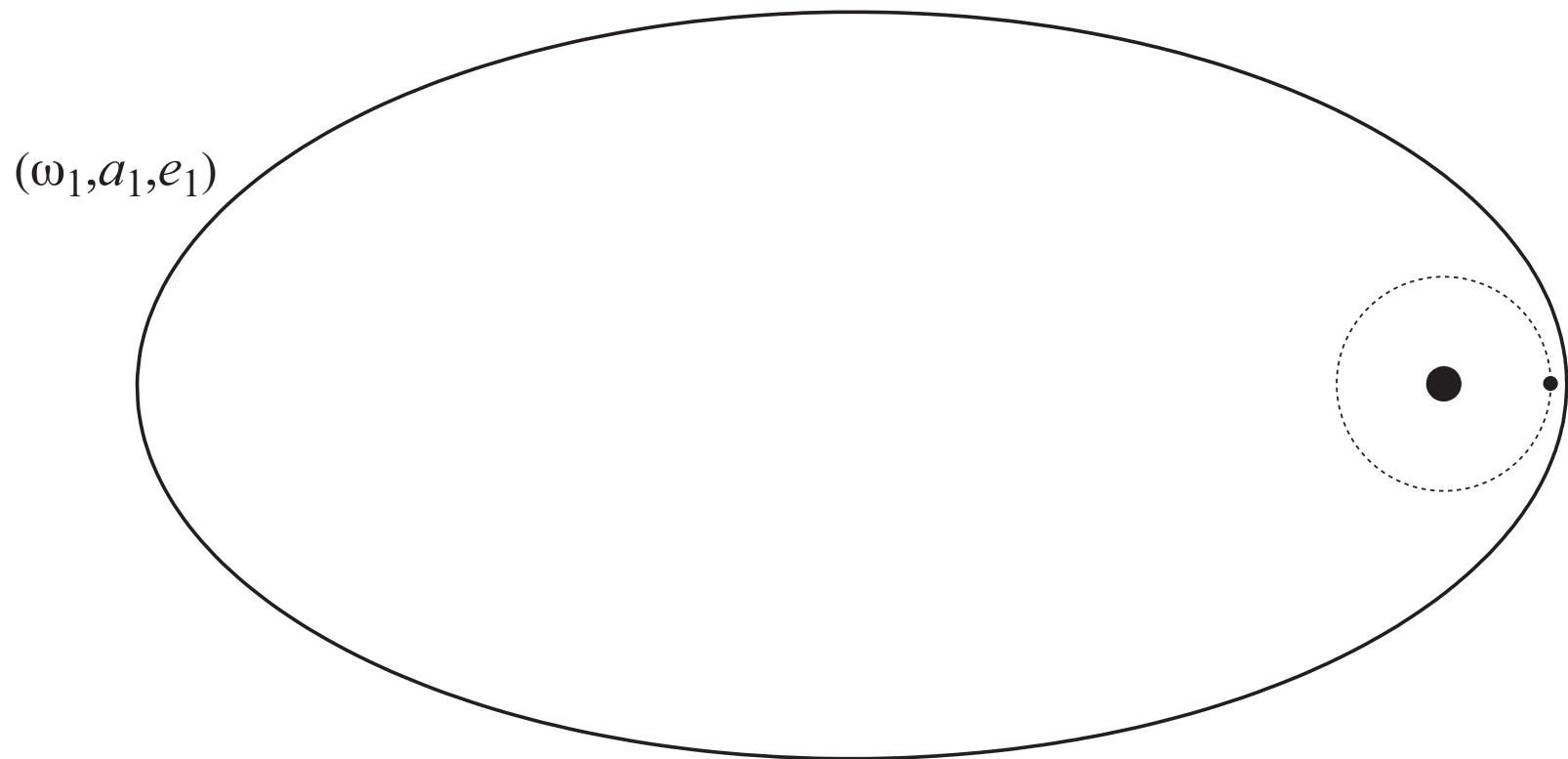
- Sensitive dependence on **argument of periapsis** ω



In rotating frame where m_1, m_2 are fixed

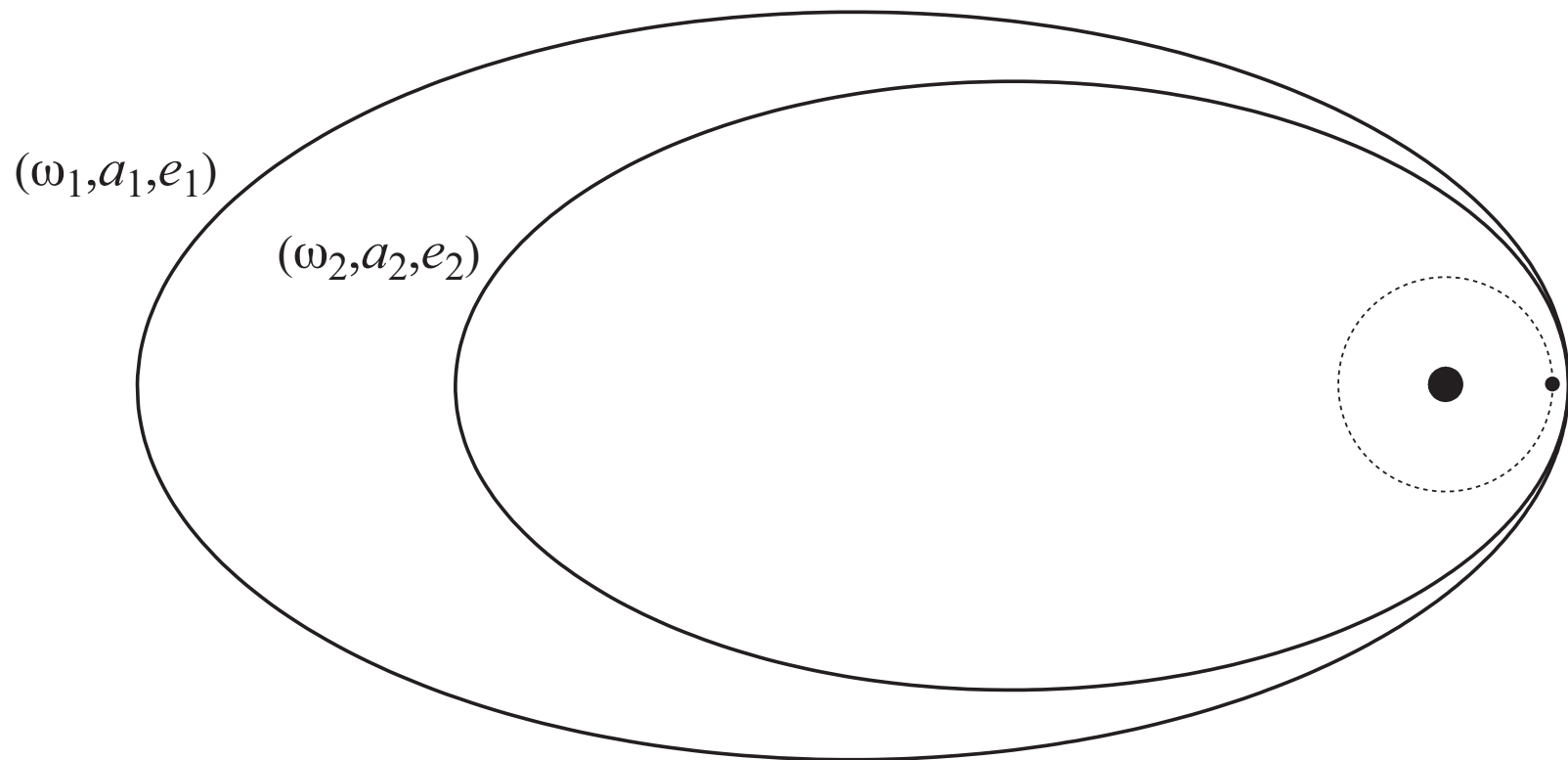
Kicks at periapsis

- Construct **update map** $(\omega_1, a_1, e_1) \mapsto (\omega_2, a_2, e_2)$ using average perturbation per orbit by smaller mass



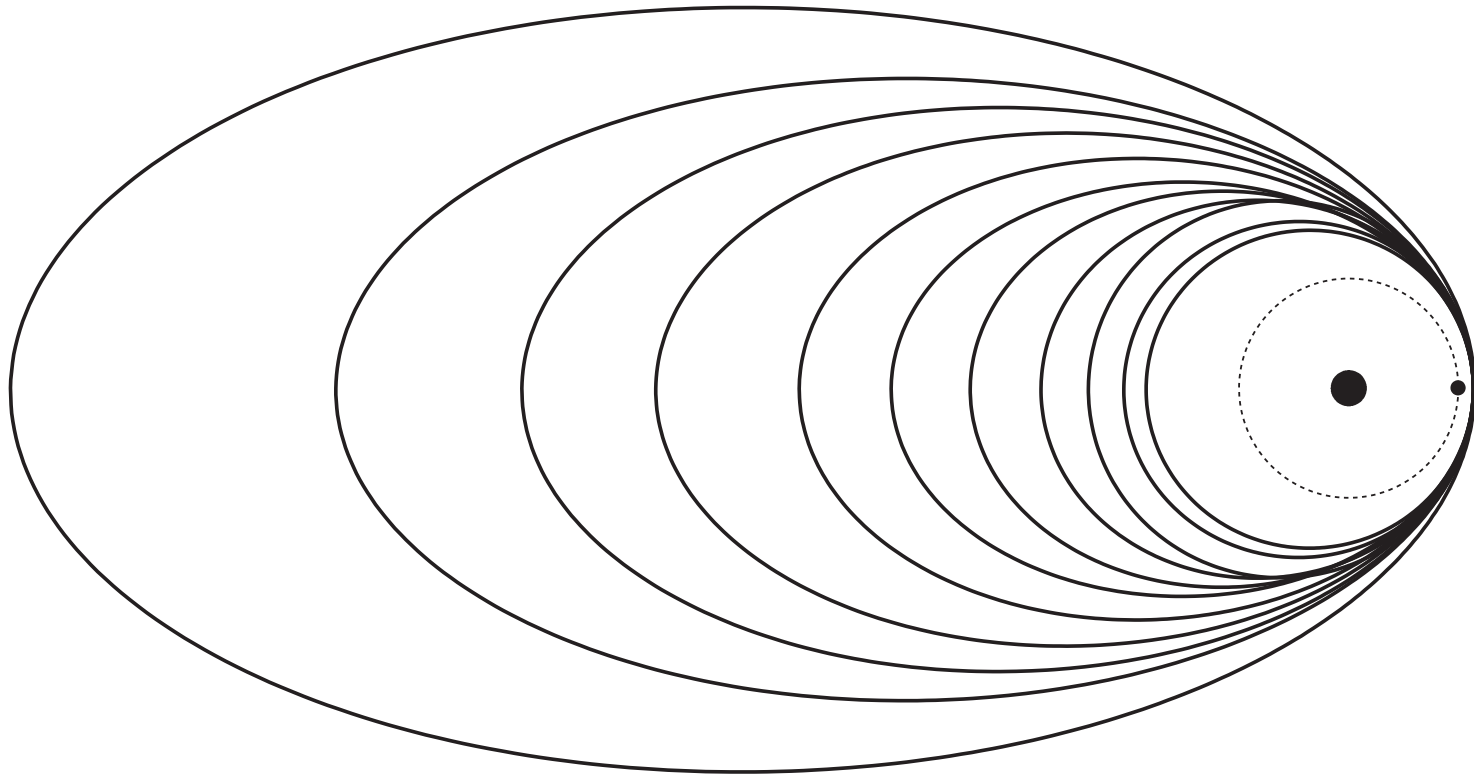
Kicks at periapsis

- Construct **update map** $(\omega_1, a_1, e_1) \mapsto (\omega_2, a_2, e_2)$ using average perturbation per orbit by smaller mass



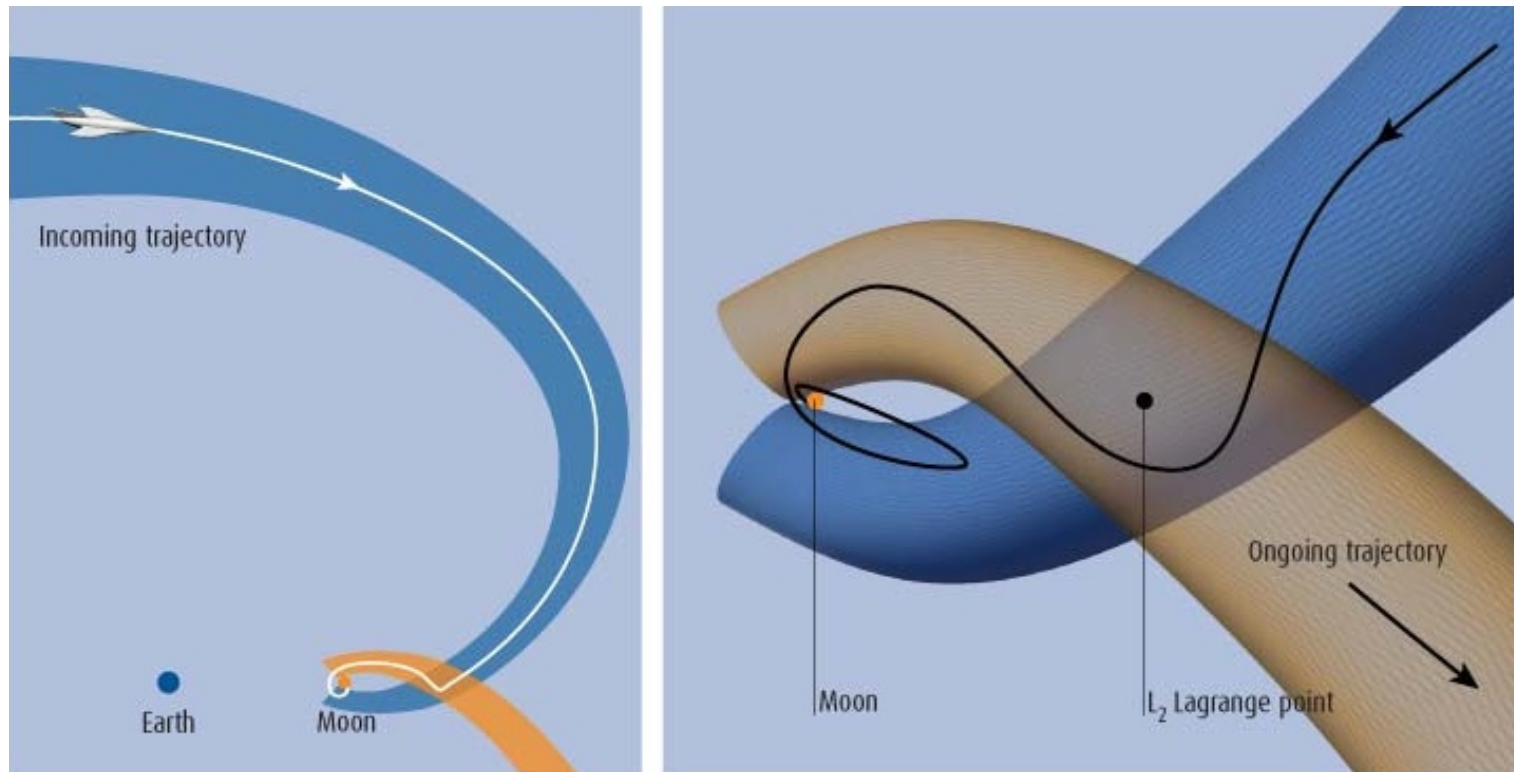
Not hyperbolic swing-by

- Occur **outside sphere of influence** (Hill radius)
 - not the close, hyperbolic swing-bys of Voyager



Capture by secondary

- Dynamically connected to capture thru tubes



Capture by secondary

- Particle assumed on **near-Keplerian orbit** around m_1
- In the frame co-rotating with m_2 and m_1 ,

$$H_{\text{rot}}(l, \omega, L, G) = K(L) + \mu R(l, \omega, L, G) - G,$$

in Delaunay variables

- Evolution is Hamilton's equations:

$$\frac{d}{dt}(l, \omega, L, G) = f(l, \omega, L, G)$$

- Jacobi constant, $C_J = -2H_{\text{rot}}$
conserved along trajectories

Change in orbital elements over one particle orbit

- Evolution of G (angular momentum)

$$\frac{dG}{dt} = -\mu \frac{\partial R}{\partial \omega},$$

- Picard's approximation:

$$\begin{aligned} \Delta G &= -\mu \int_{-T/2}^{T/2} \frac{\partial R}{\partial \omega} dt \\ &= -\frac{\mu}{G} \left[\left(\int_{-\pi}^{\pi} \left(\frac{r}{r_2} \right)^3 \sin(\omega + \nu - t(\nu)) d\nu \right) - \sin \omega \left(2 \int_0^{\pi} \cos(\nu - t(\nu)) d\nu \right) \right] \end{aligned}$$

- $\Delta K =$ **Keplerian energy change** over an orbit

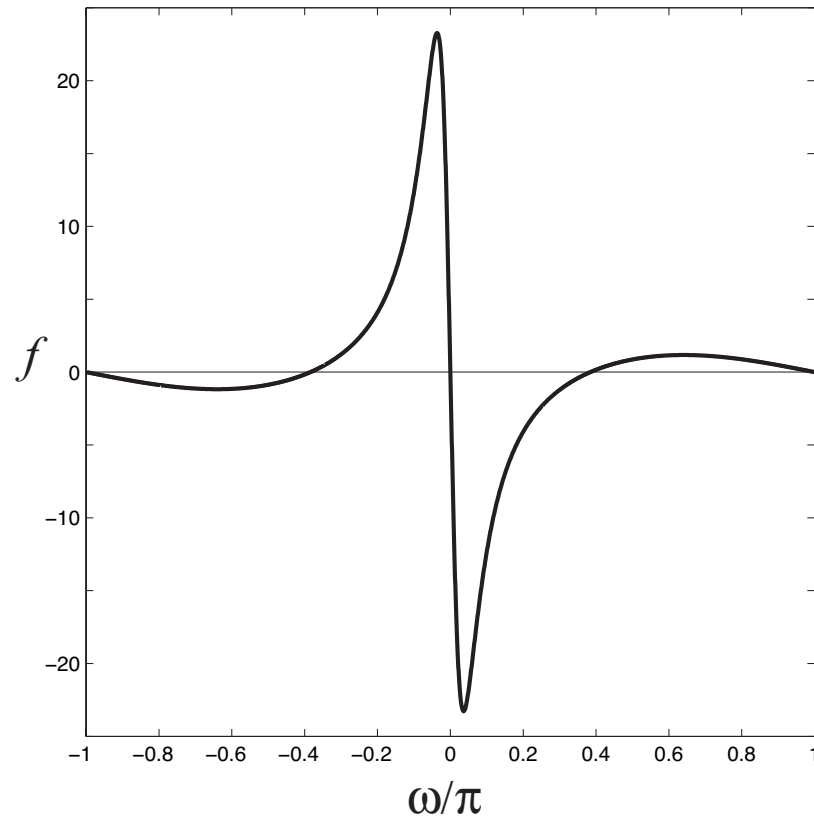
$$\Delta K = \Delta G - \mu \Delta R$$

Energy kick function

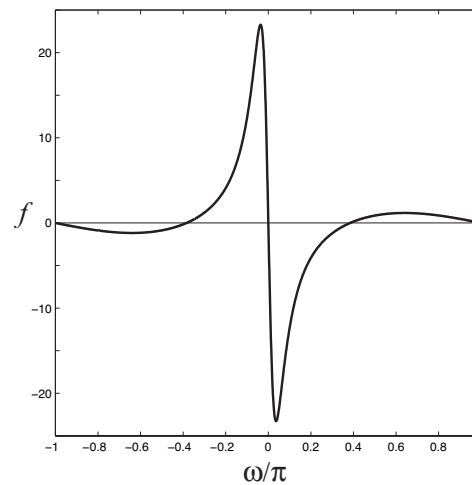
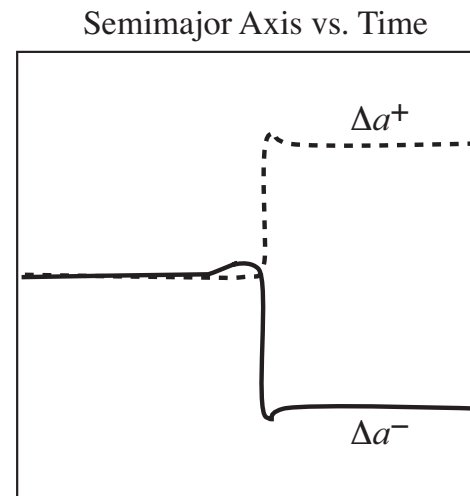
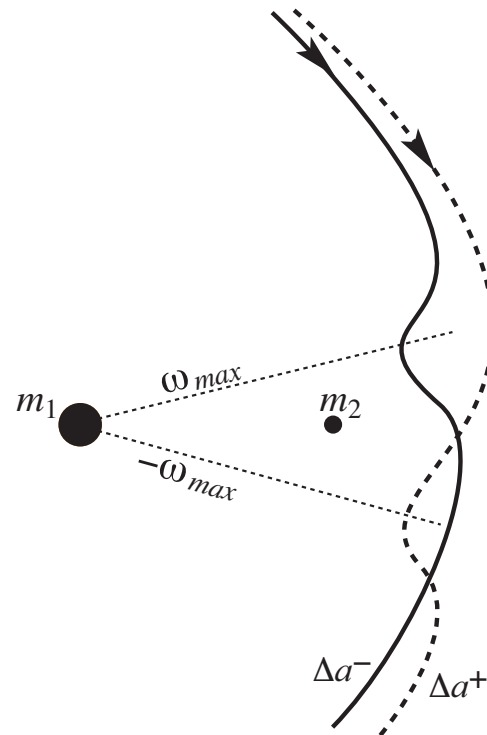
□ Changes have form

$$\Delta K = \mu f(\omega),$$

f is the **energy kick function** with parameters K, C_J



Maximum changes on either side of perturber



The periapsis kick map (Keplerian Map)

□ Cumulative effect of **consecutive passes** by perturber

□ Can construct an **update map**

$(\omega_{n+1}, K_{n+1}) = F(\omega_n, K_n)$ on the cylinder $\Sigma = S^1 \times \mathbb{R}$,
i.e., $F : \Sigma \rightarrow \Sigma$ where

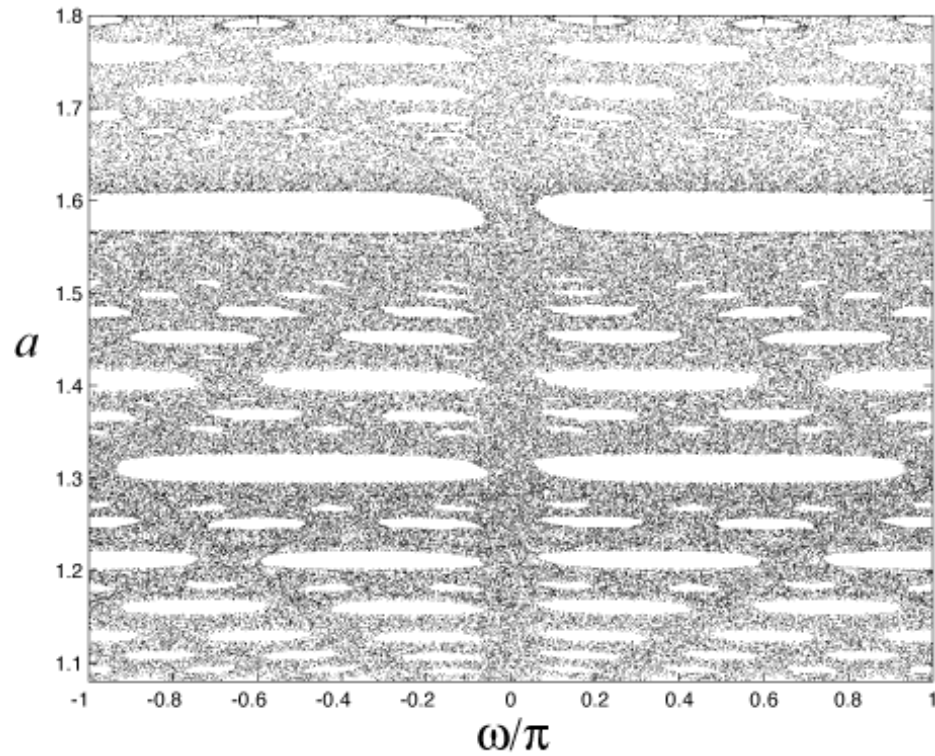
$$\begin{pmatrix} \omega_{n+1} \\ K_{n+1} \end{pmatrix} = \begin{pmatrix} \omega_n - 2\pi(-2(K_n + \mu f(\omega_n)))^{-3/2} \\ K_n + \mu f(\omega_n) \end{pmatrix}$$

□ **Area-preserving (symplectic twist) map**

□ Example: particle in Jupiter-Callisto system

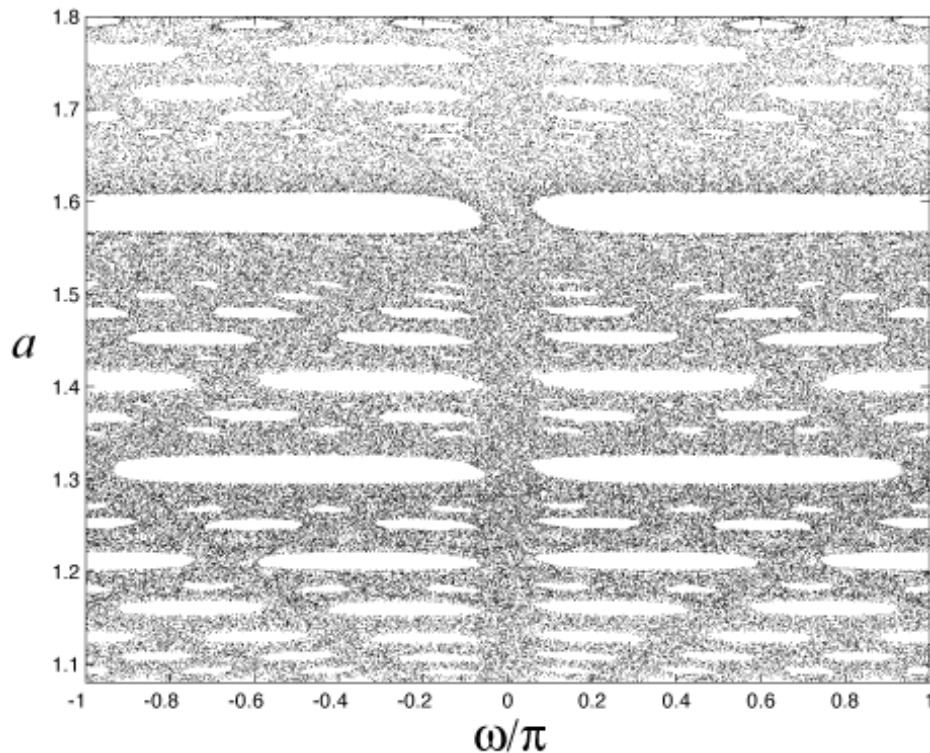
$$\mu = 5 \times 10^{-5}$$

Verification of Keplerian map: phase portrait

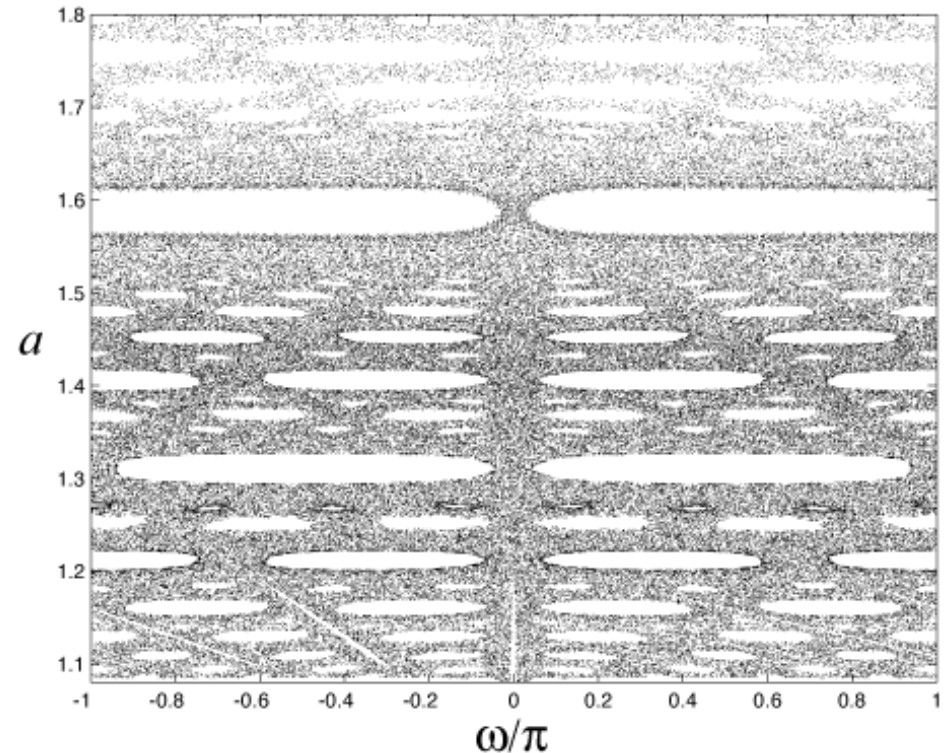


Keplerian map

Verification of Keplerian map: phase portrait



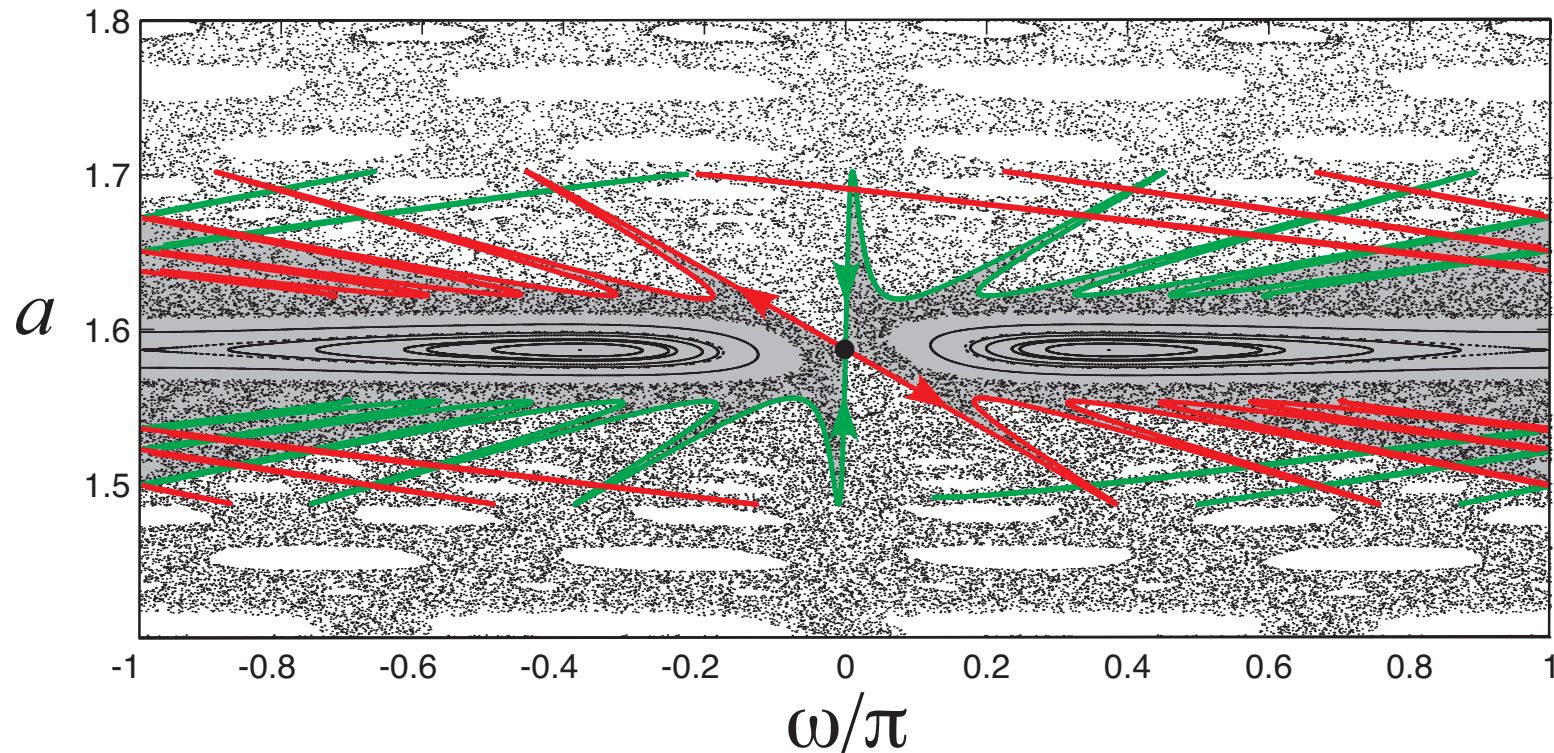
Keplerian map



numerical integration of ODEs

- Keplerian map = fast orbit propagator
- preserves phase space features
 - but breaks left-right symmetry present in original system
 - can be removed using another method (Hamilton-Jacobi)

Dynamics of Keplerian map

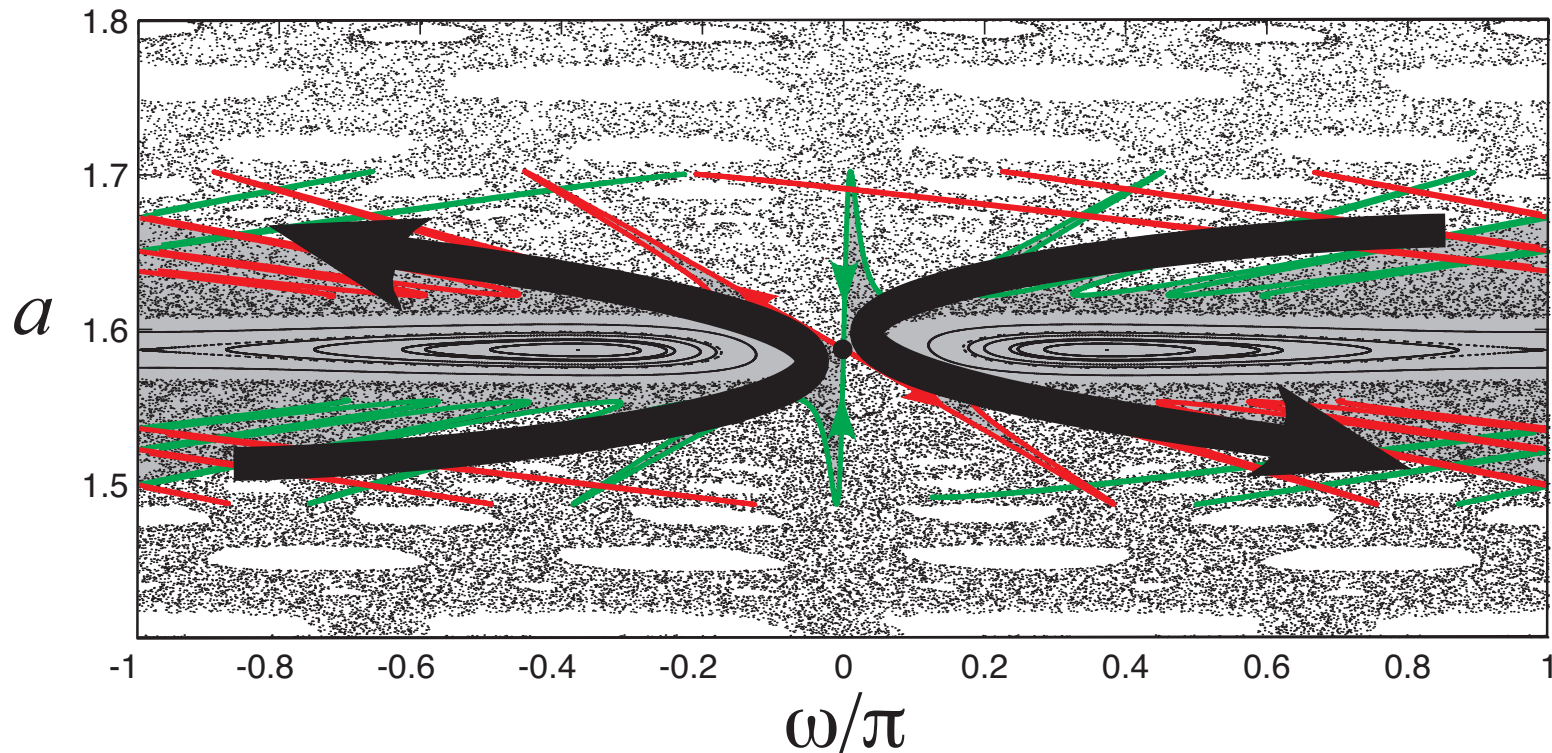


Resonance zone³

□ Structured motion around resonance zones

³in the terminology of MacKay, Meiss, and Percival [1987]

Dynamics of Keplerian map



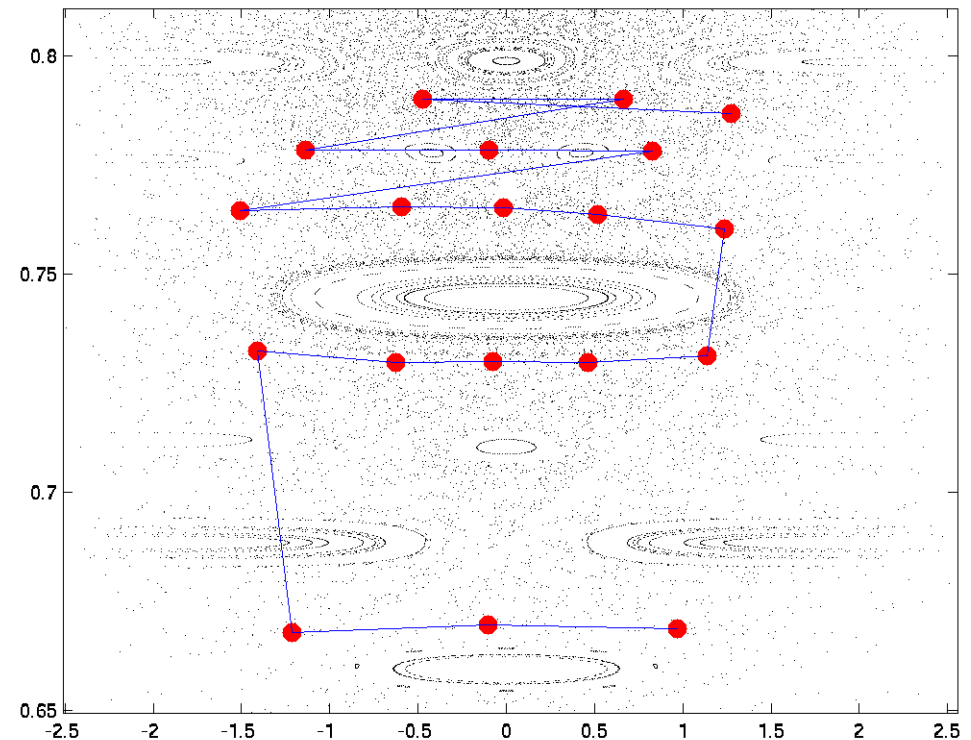
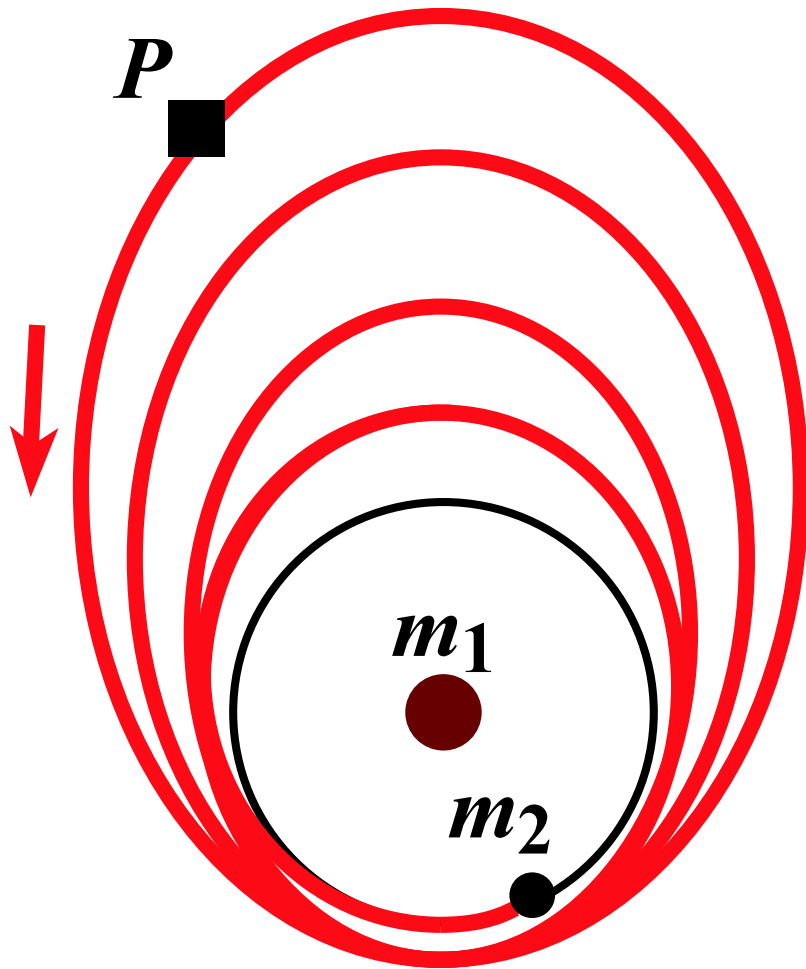
Resonance zone⁴

□ Structured motion around resonance zones

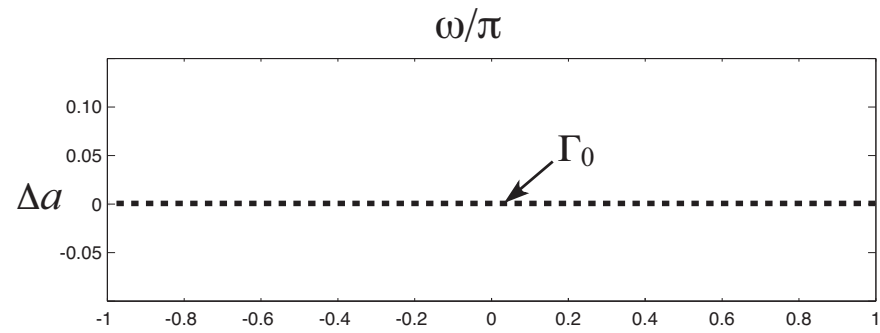
⁴in the terminology of MacKay, Meiss, and Percival [1987]

Large orbit changes via multiple resonance zones

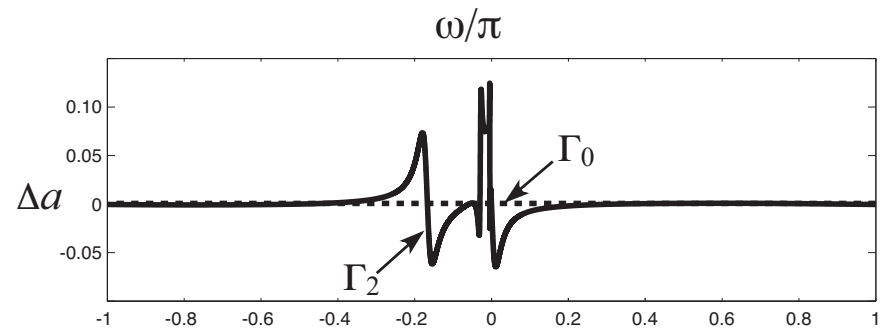
- multiple flybys for orbit reduction or expansion



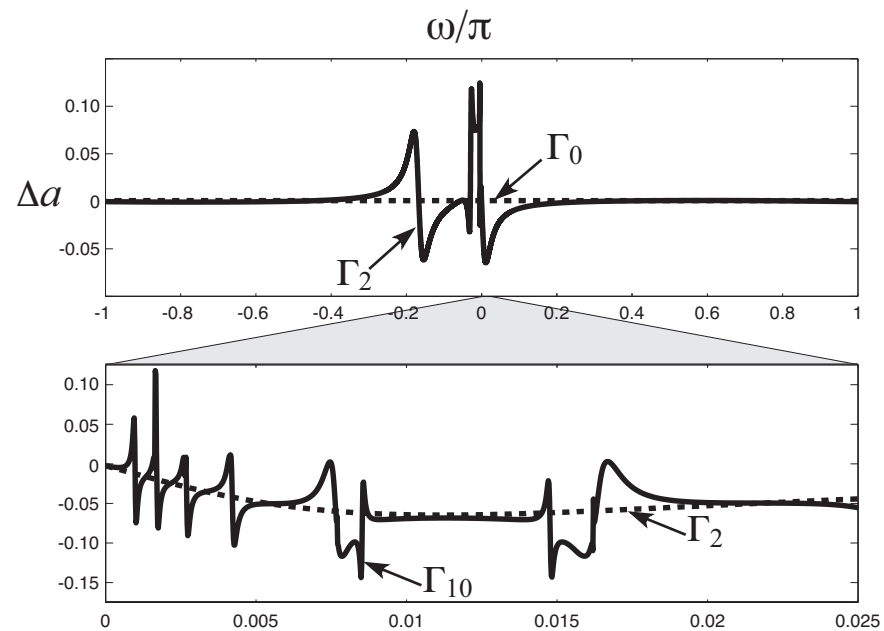
Large orbit changes, $\Gamma_n = F^n(\Gamma_0)$



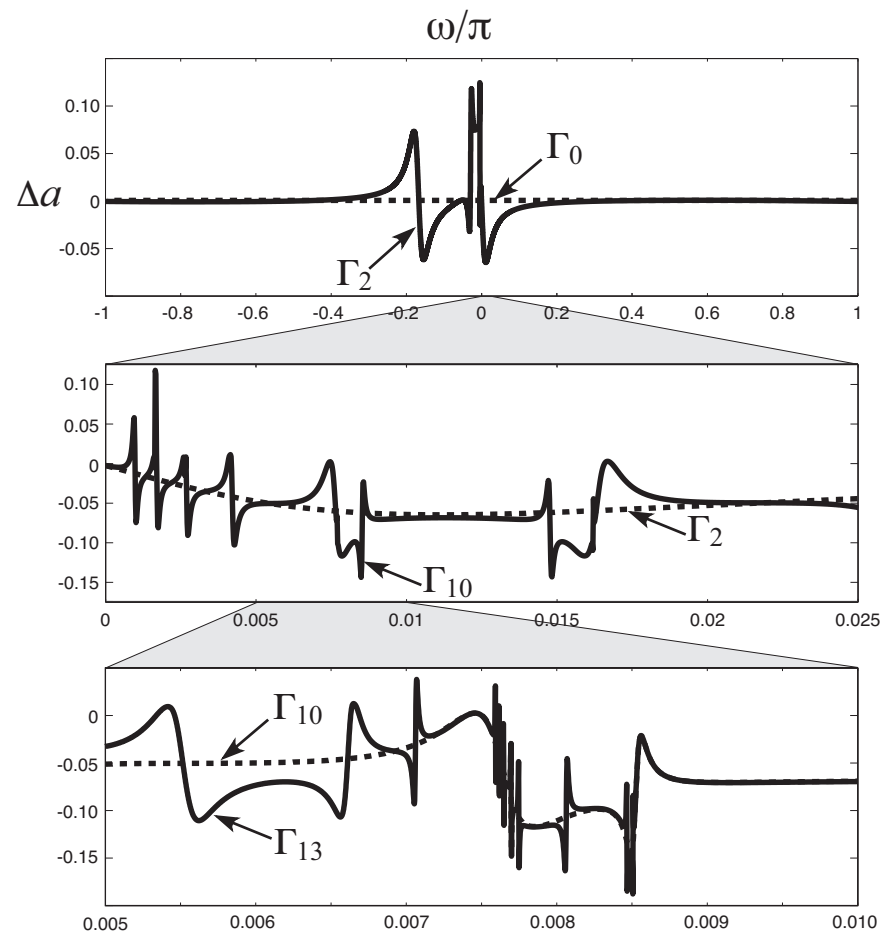
Large orbit changes, $\Gamma_n = F^n(\Gamma_0)$



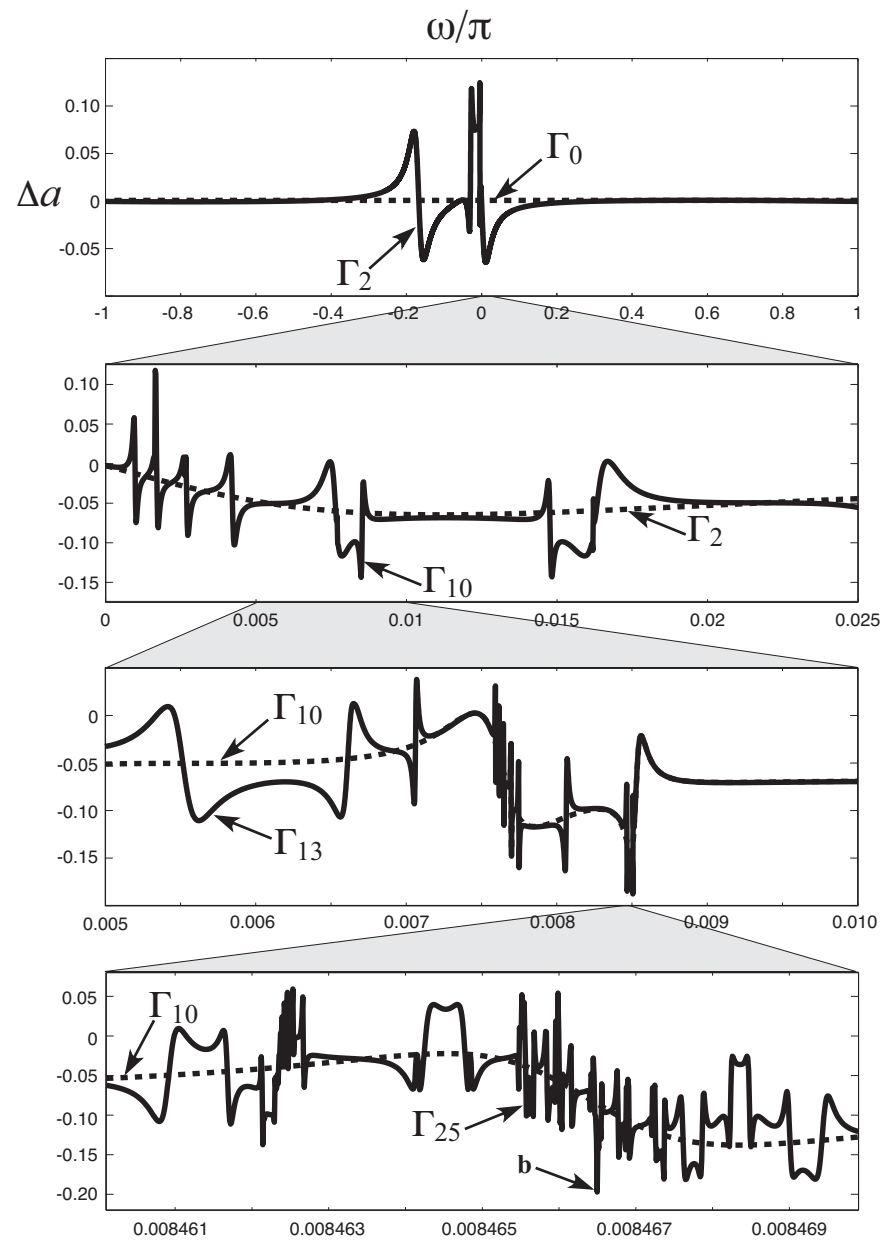
Large orbit changes, $\Gamma_n = F^n(\Gamma_0)$



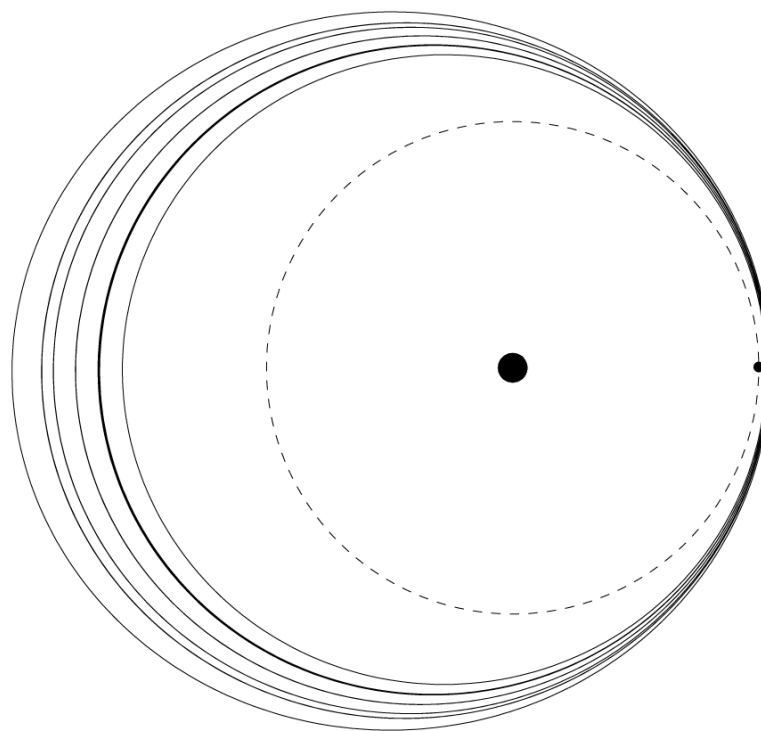
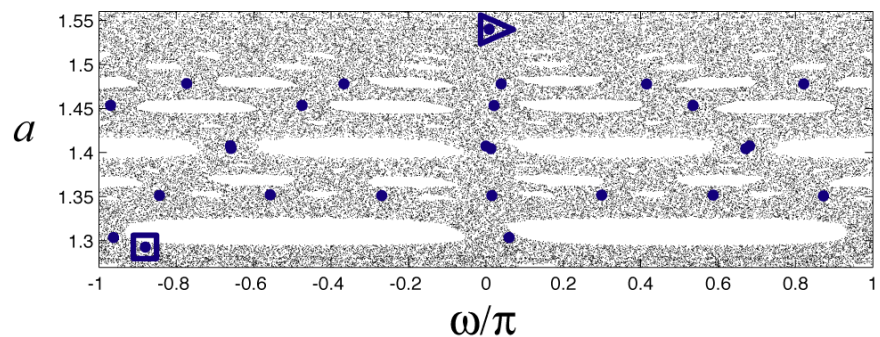
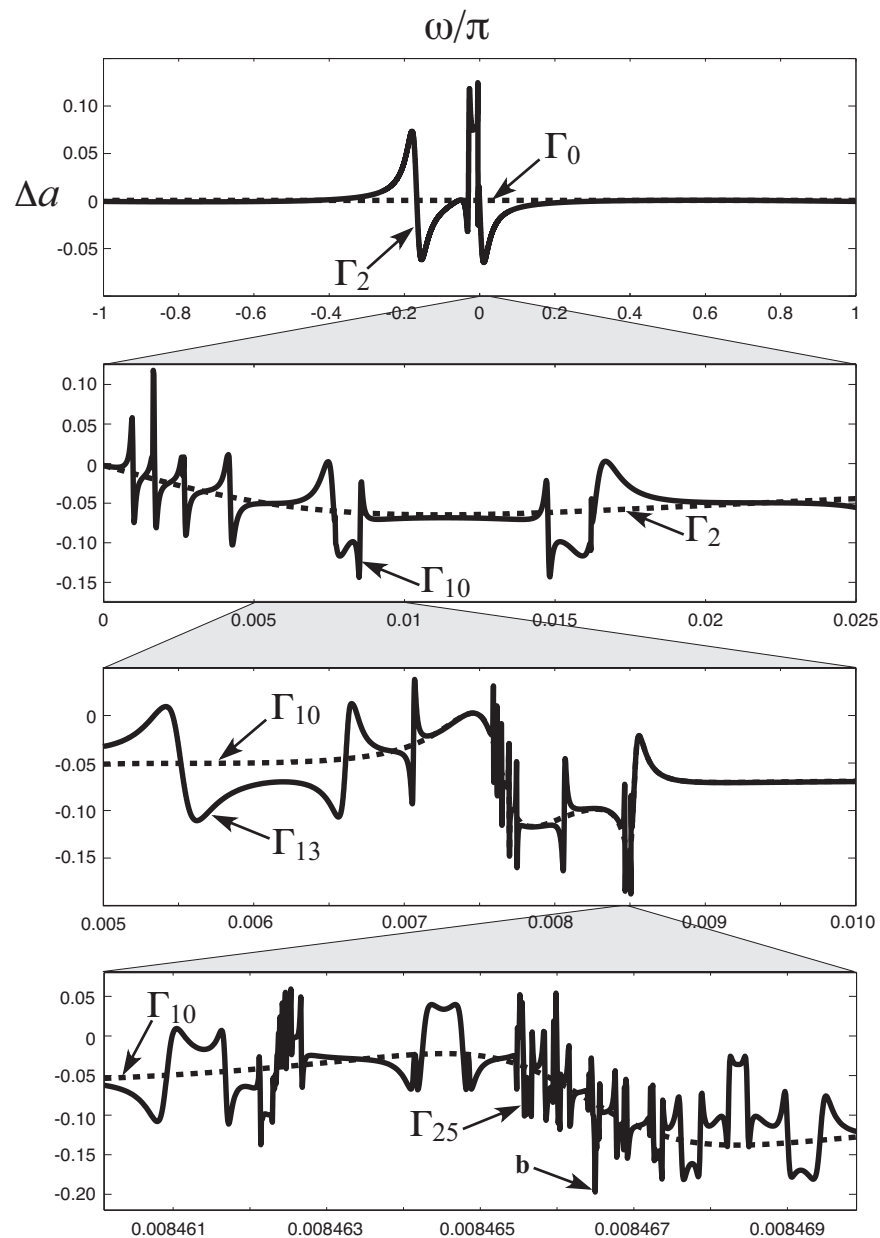
Large orbit changes, $\Gamma_n = F^n(\Gamma_0)$



Large orbit changes, $\Gamma_n = F^n(\Gamma_0)$

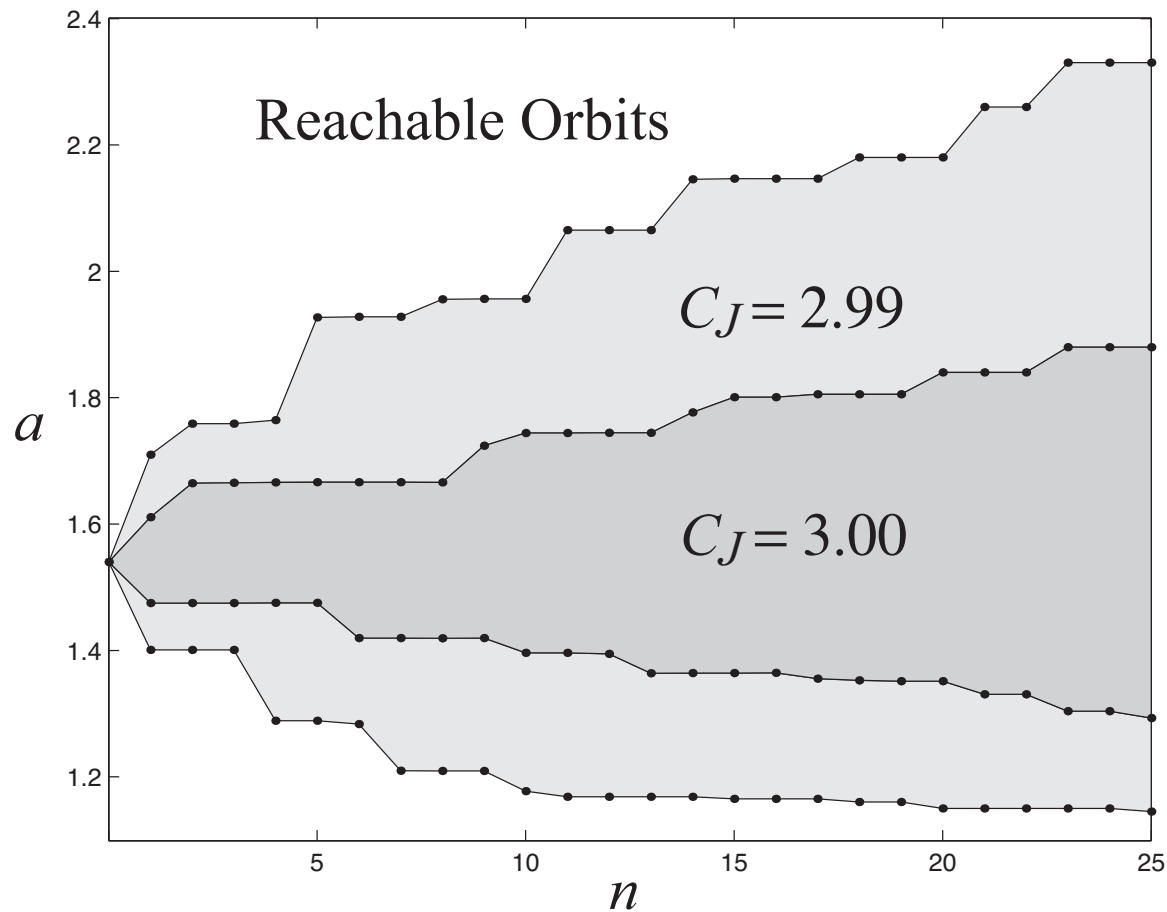


Large orbit changes, $\Gamma_n = F^n(\Gamma_0)$



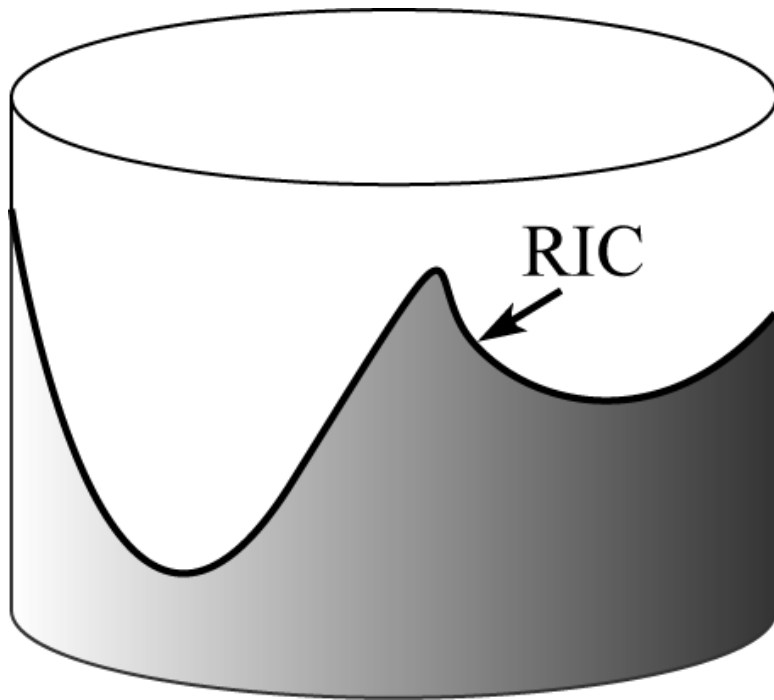
example trajectory

Reachable orbits and diffusion

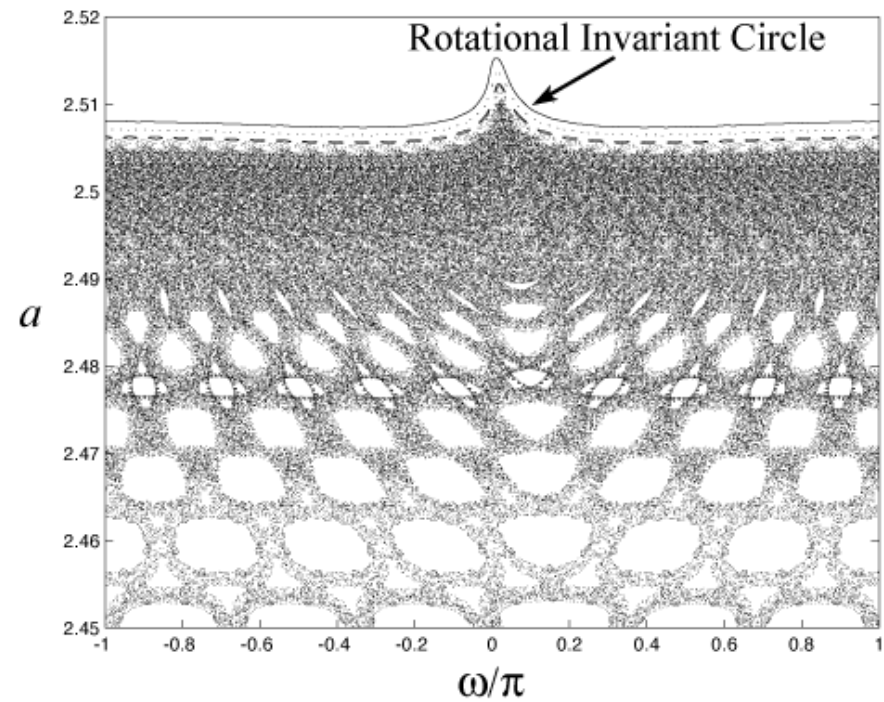


- Diffusion in semimajor axis
- ... increases as C_J decreases (larger kicks)

Reachable orbits: upper boundary for small μ

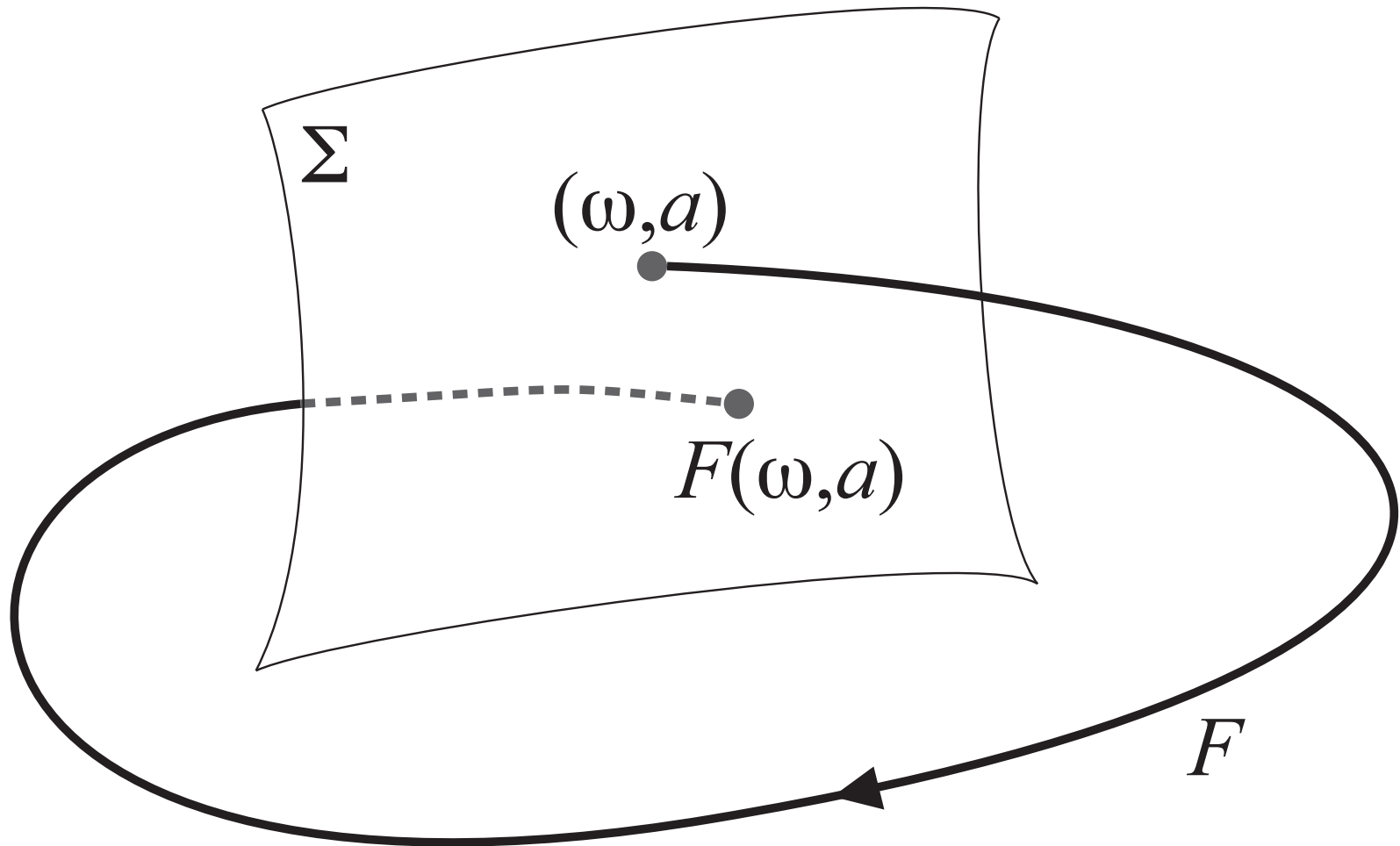


A rotational invariant circle (RIC)



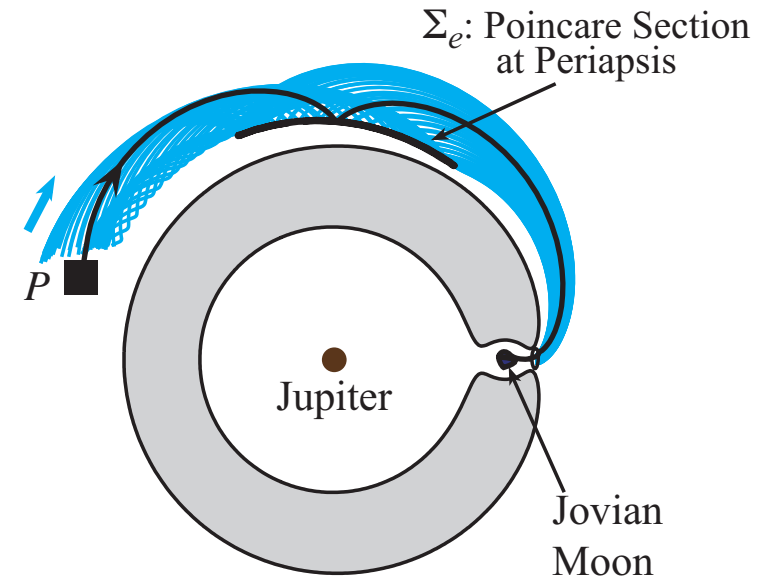
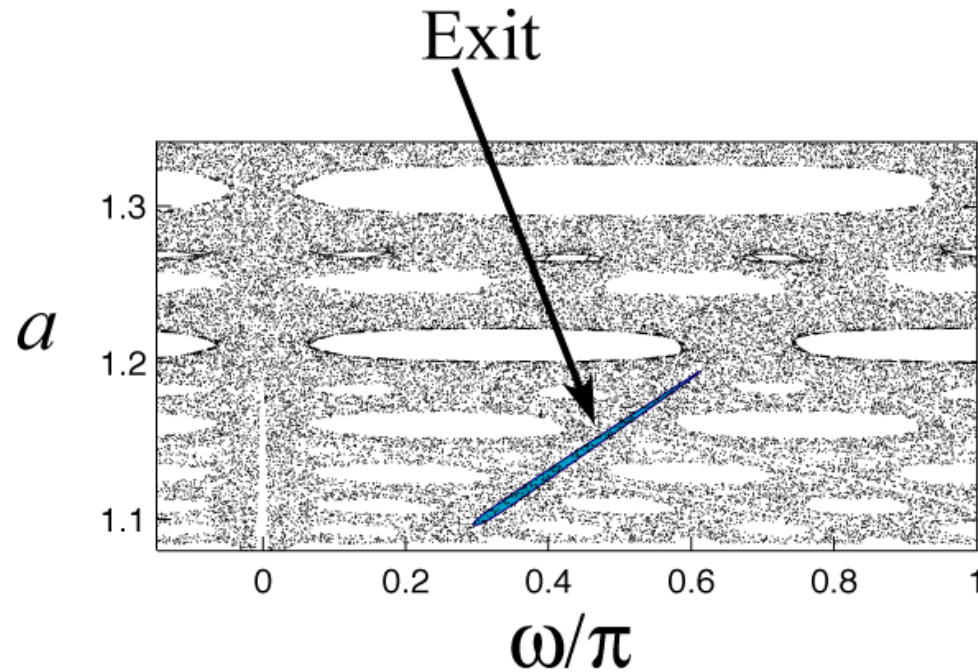
RIC found in Keplerian map for $\mu = 5 \times 10^{-6}$

Identify Keplerian map as Poincaré return map



- **Poincaré map at periapsis** in orbital element space
- $F : \Sigma \rightarrow \Sigma$ where $\Sigma = \{l = 0 \mid C_J = \text{constant}\}$

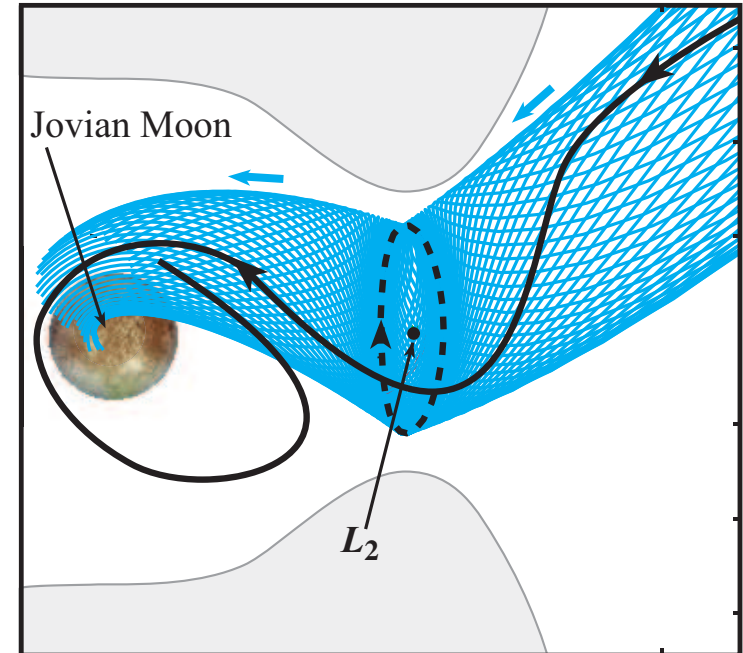
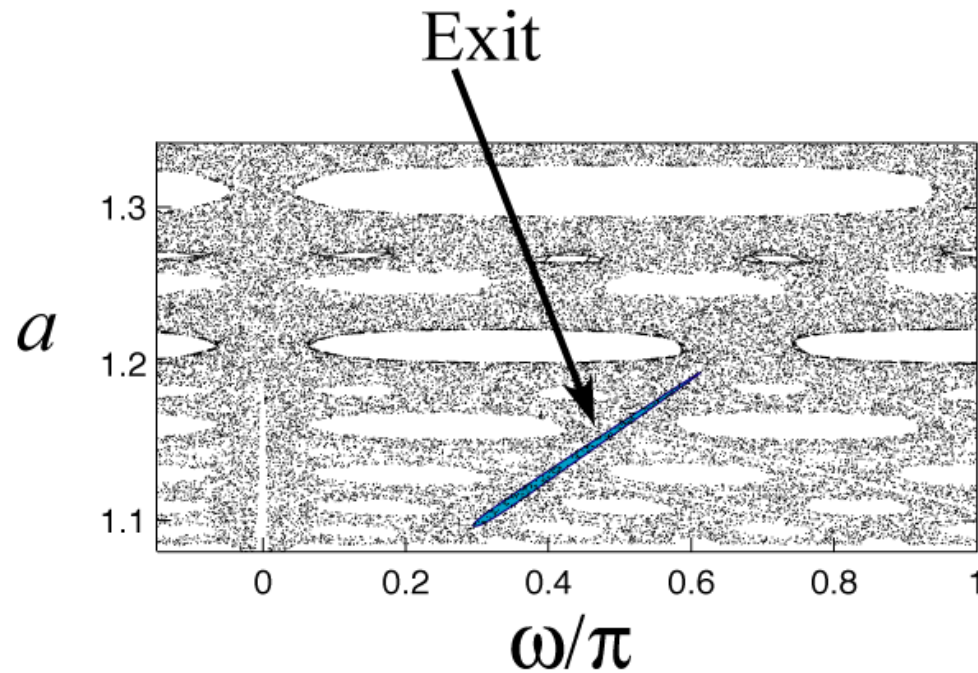
Relationship to capture around perturber



exit from jovicentric to moon region

□ **Exit**: where tube of capture orbits intersects Σ

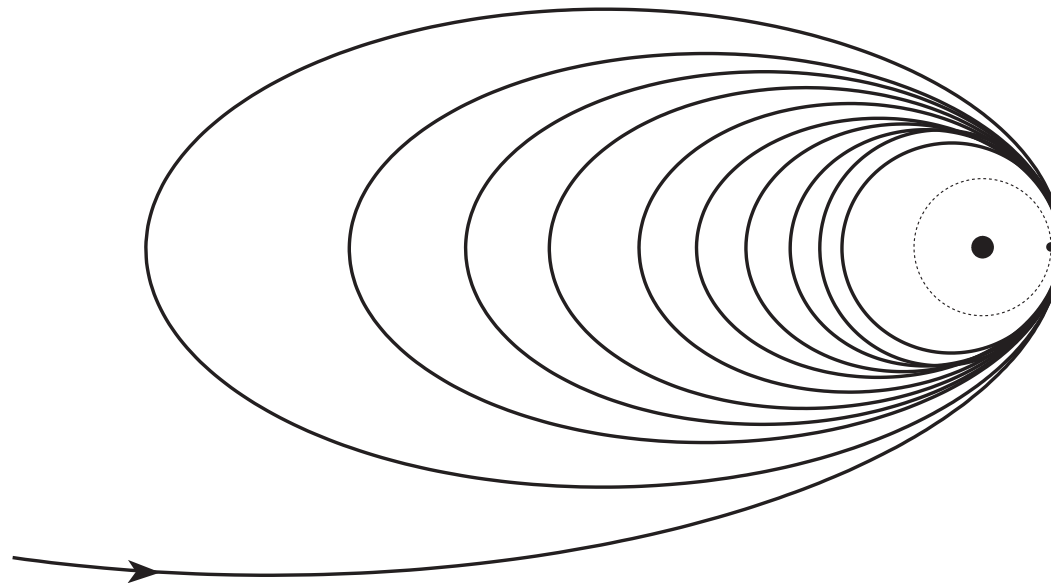
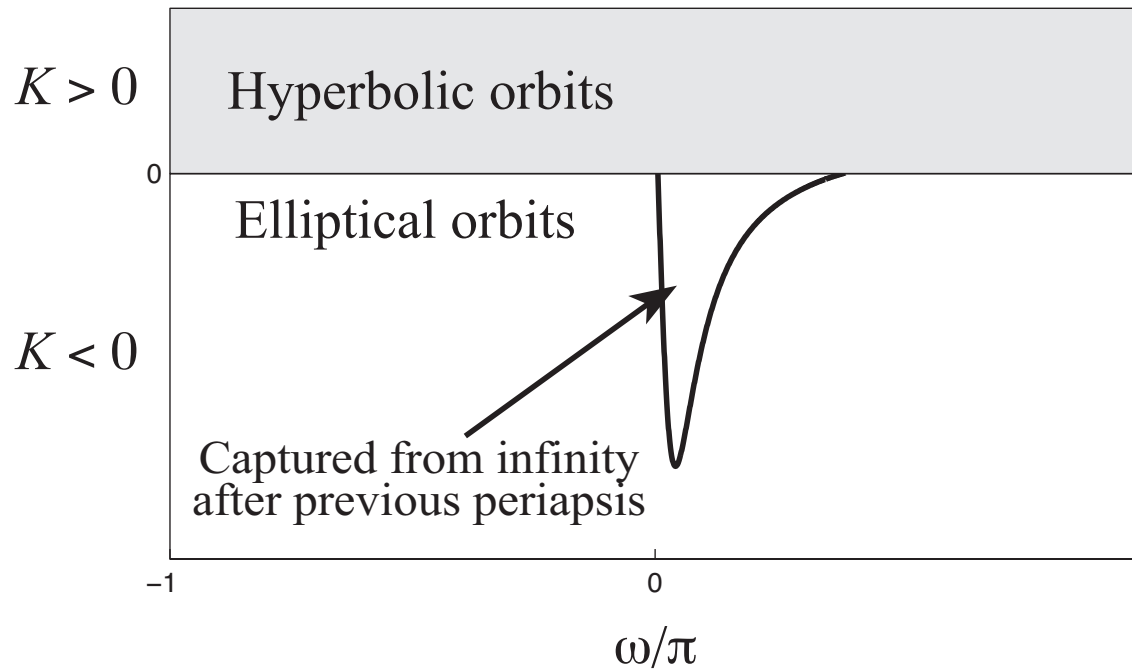
Relationship to capture around perturber



exit from jovicentric to moon region

- **Exit**: where tube of capture orbits intersects Σ
- Orbits reaching exit are **ballistically captured**, passing by L_2

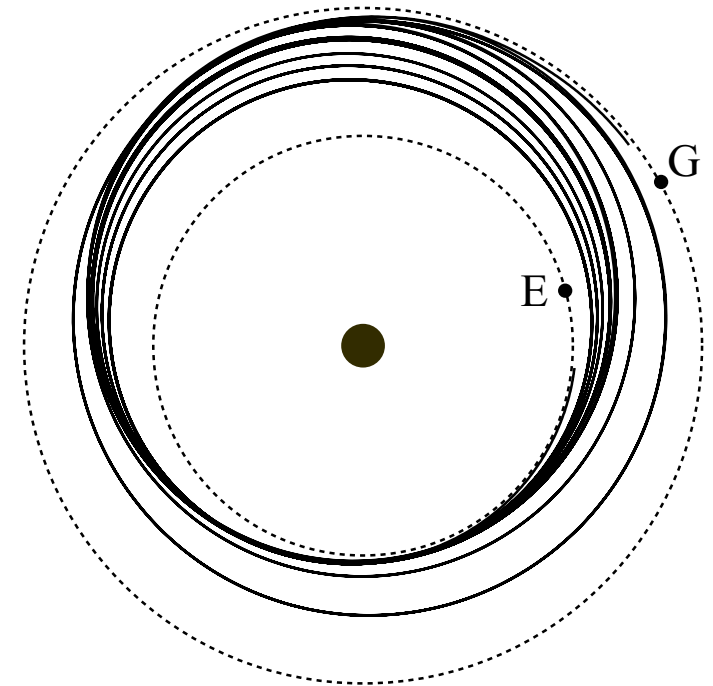
Relationship to capture from infinity



Final word about Keplerian map

□ Extensions:

- out of plane motion (**4D map**)
- control in the presence of uncertainty
- eccentric orbits for the perturbers
- multiple perturbers
transfer from one body to another



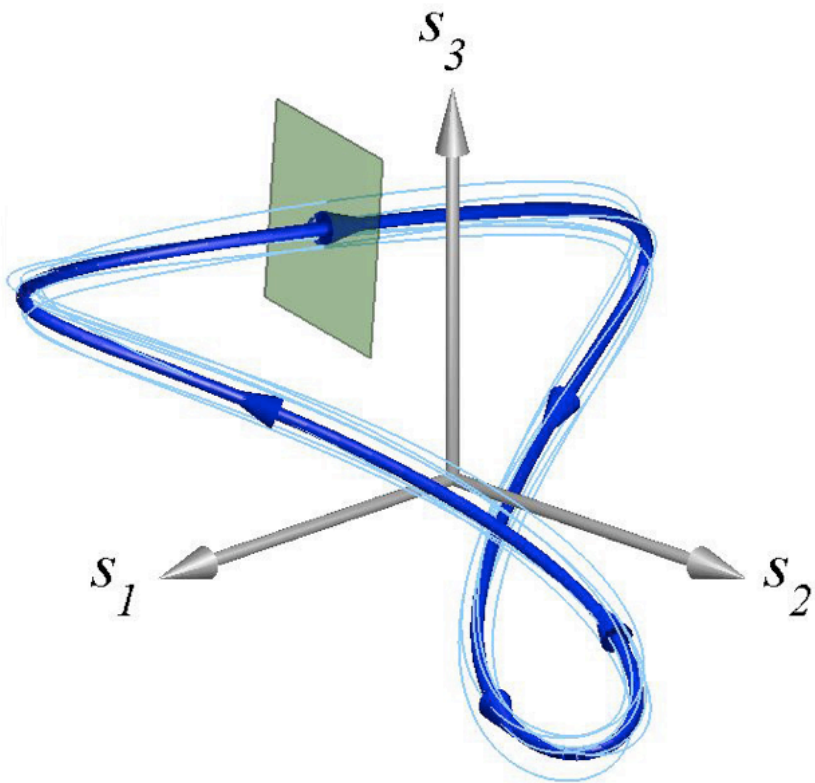
- Consider other problems with localized perturbations?
 - chemistry, vortex dynamics, ...

Reference:

Ross & Scheeres, *SIAM J. Applied Dynamical Systems*, 2007.

Separatrices: biomechanics

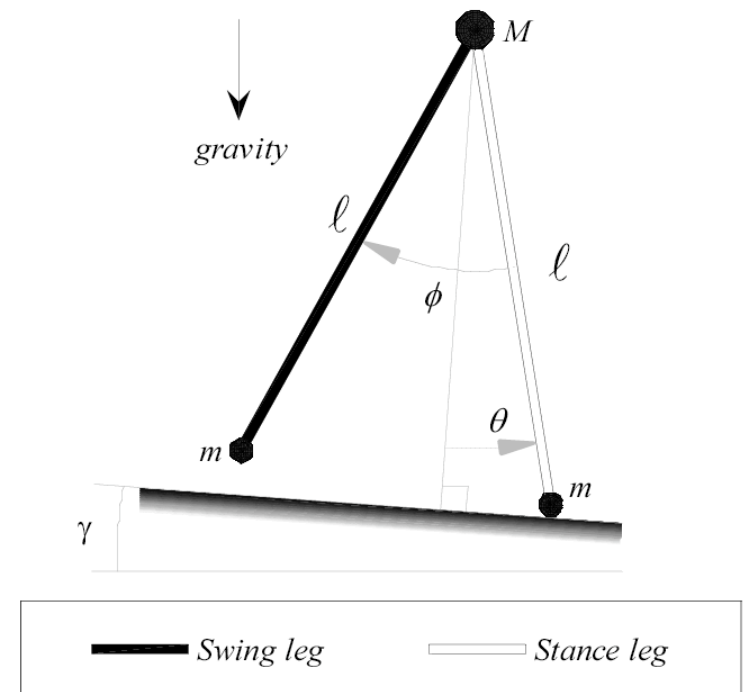
- **Boundaries between qualitatively (functionally) different kinds of behavior**
- **For example, walking or standing versus falling**
- **Based on analytical models or experimentally observed data**



Planar 2-dof model of biped walking

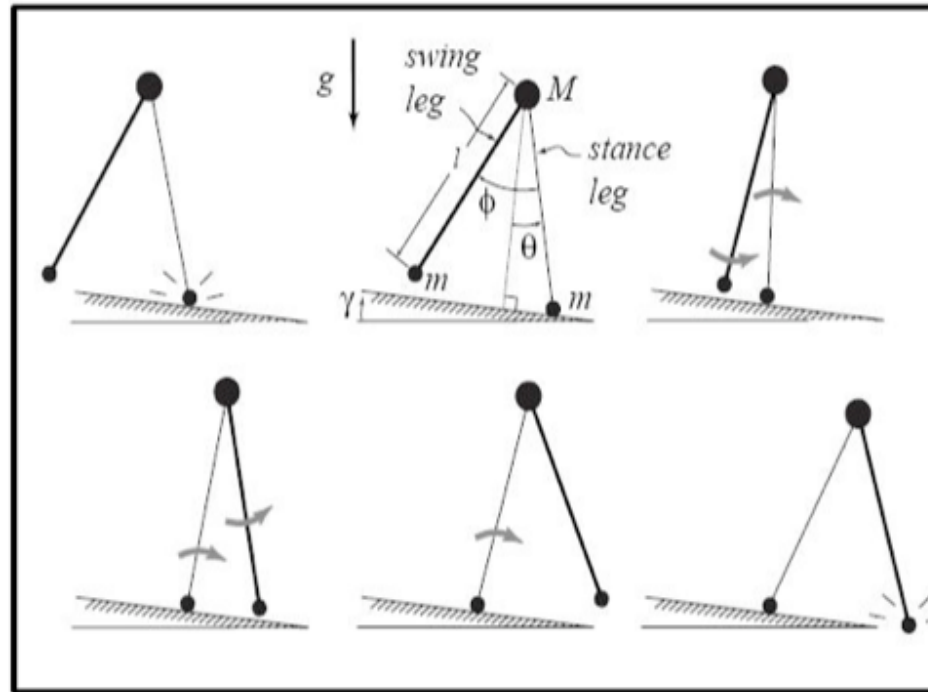
Two segment “compass biped” walker¹

- Simplest model of walking
- Double pendulum w/pin-joint at stance foot
- Point masses at hip and feet
- Massless legs of equal length
- Walks down slope
- Piecewise holonomic system
- Swing phase: Hamiltonian dynamics
- Foot-strike event: discrete, dissipative



¹McGeer 1990. Garcia et al. 1998. Norris, Marsh, Granata, Ross, 2008, Physica D.

Planar 2-dof model of biped walking

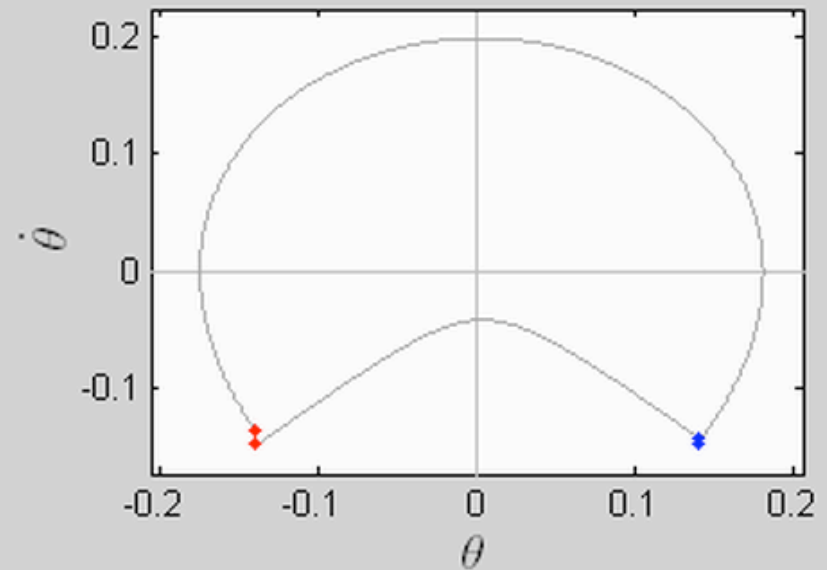
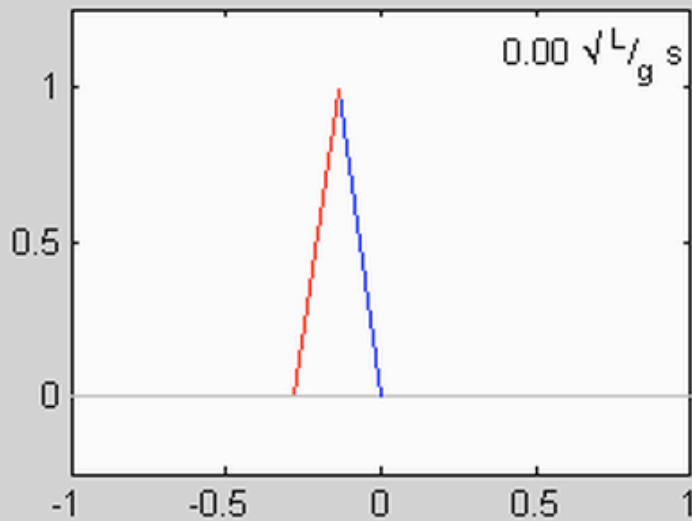
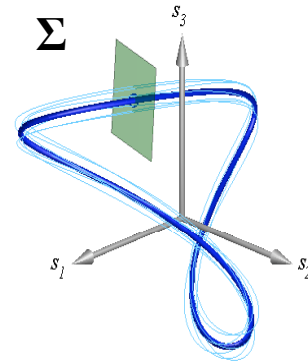


¹McGeer 1990. Garcia et al. 1998. Norris, Marsh, Granata, Ross, 2008, Physica D.

Walking solutions

Poincare section at foot-strike

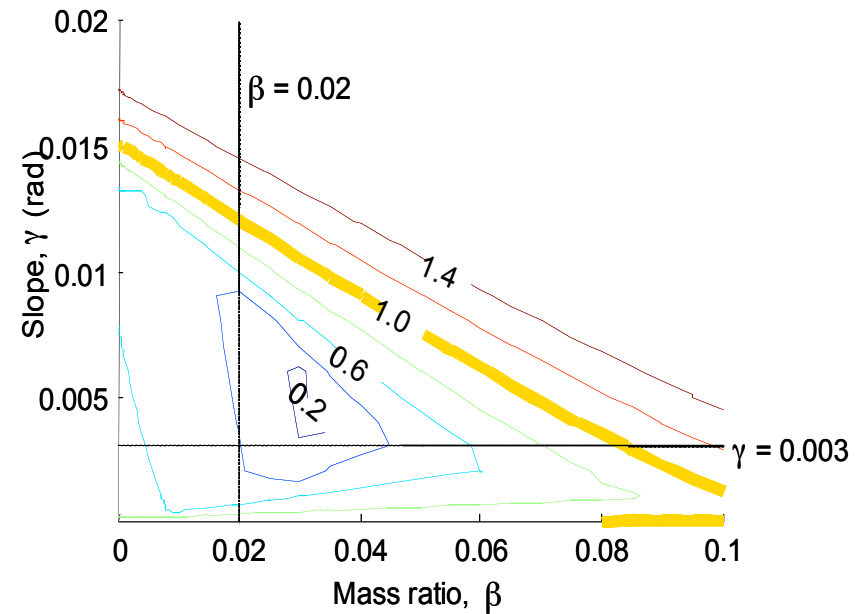
- Search for period n solutions
- Period one solution shown



Walking solution stability

Orbital stability

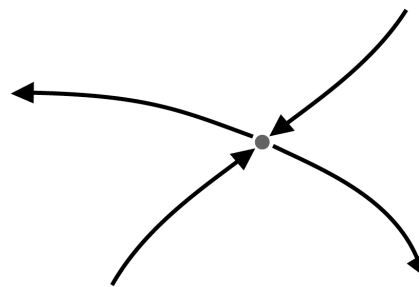
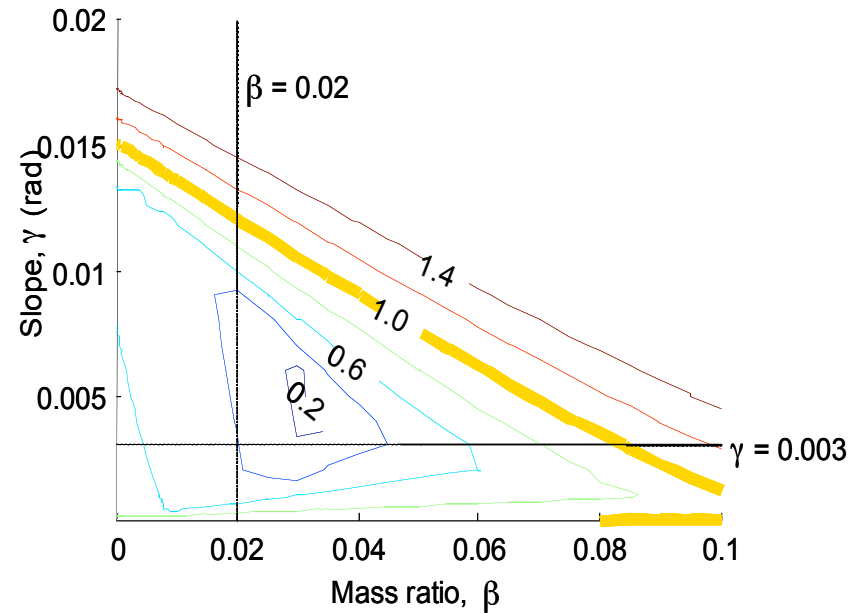
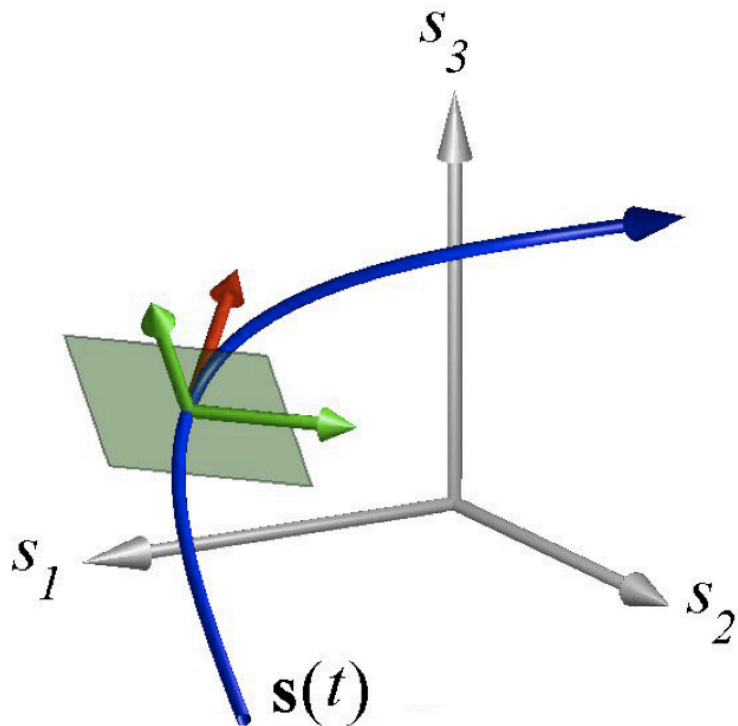
- Maximum Floquet multipliers shown



Walking solution stability

Orbital stability

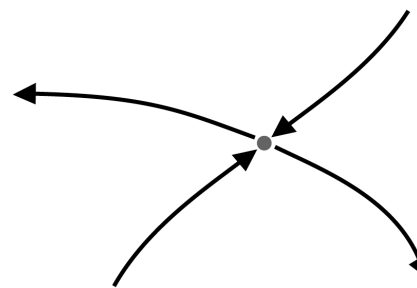
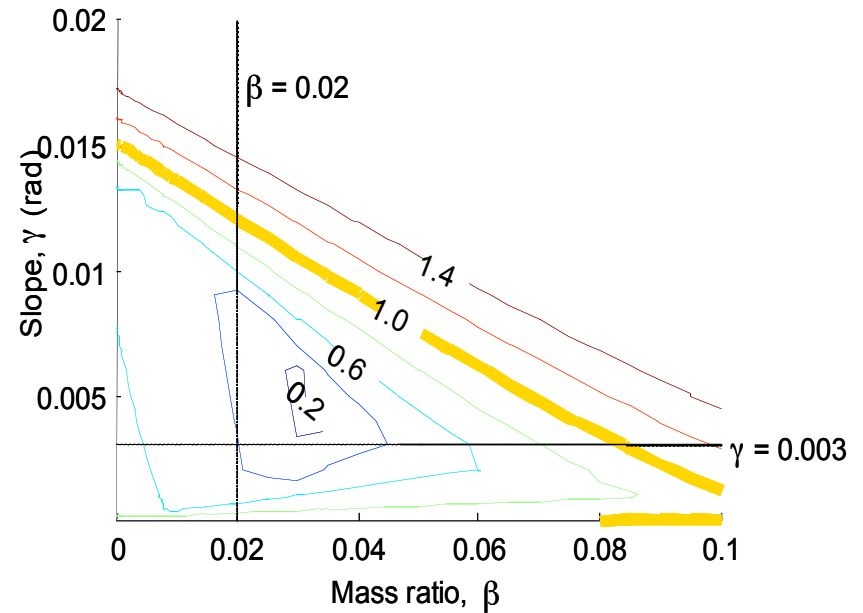
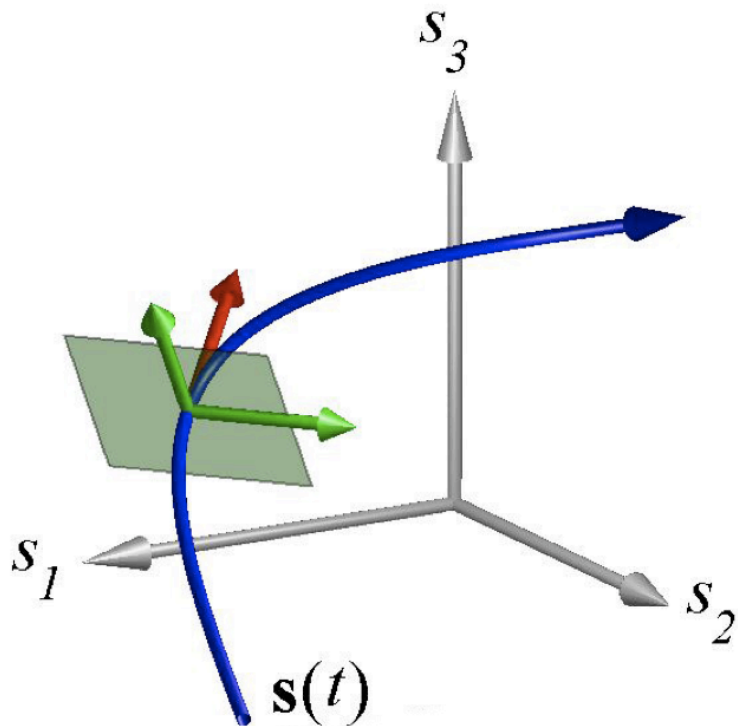
- Maximum Floquet multipliers shown
- Swing phase has one dimensional unstable direction



Walking solution stability

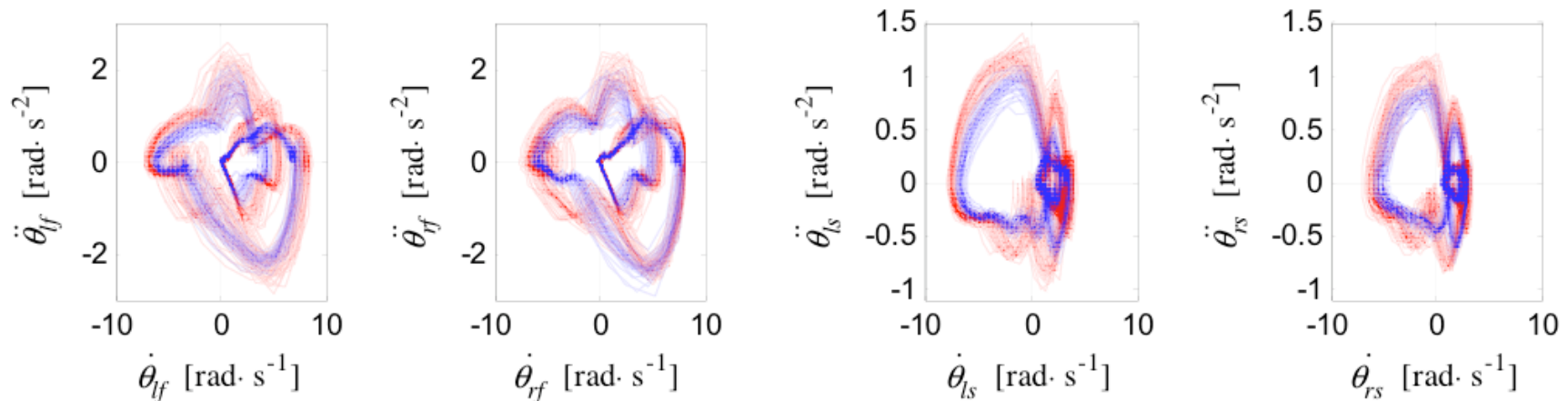
Orbital stability

- Maximum Floquet multipliers shown
- Swing phase has one dimensional unstable direction



- Dissipation at foot-strike can lead to overall stability

Analysis of bipedal walking

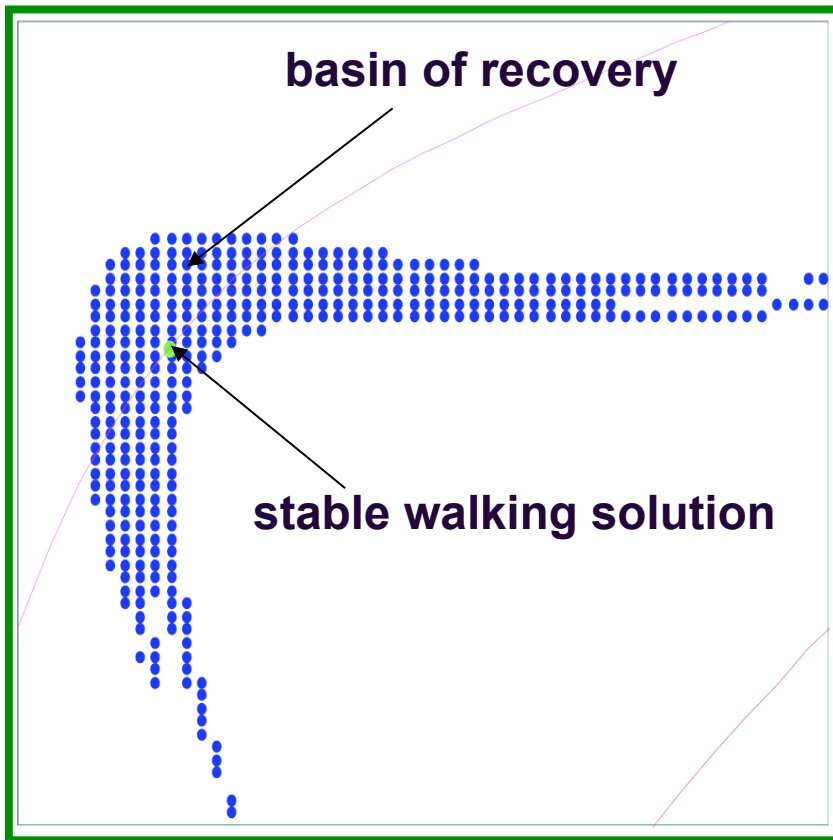


- **Walking gait data -- nearly cyclic in**
- **high dimensional phase space**

- **Orbital stability (maximum Floquet**
- **multiplier) typically around 0.7**

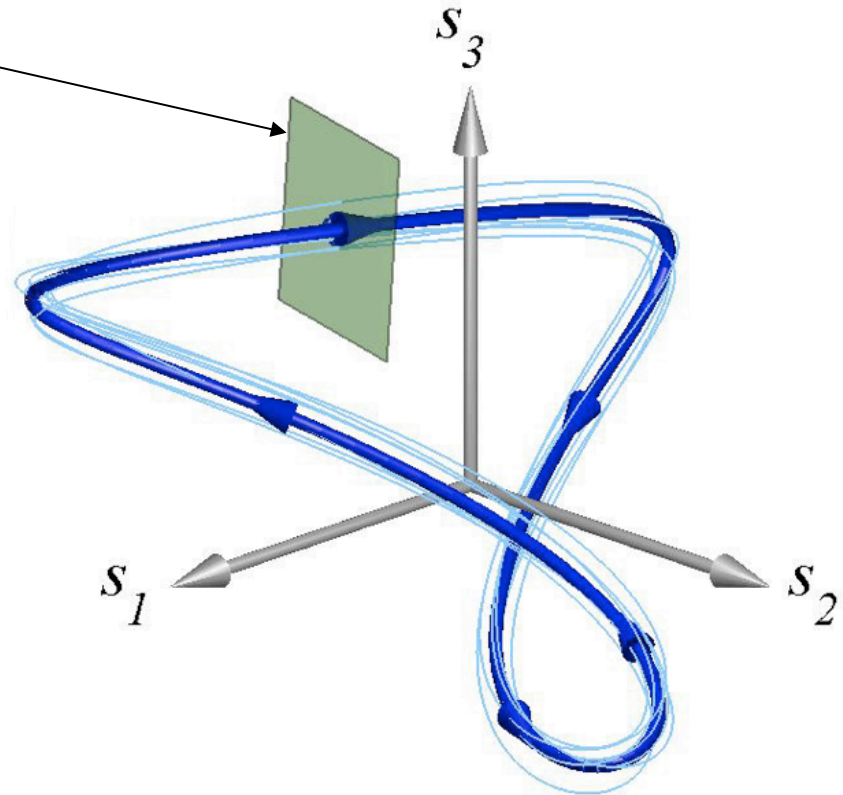


Walking solution stability



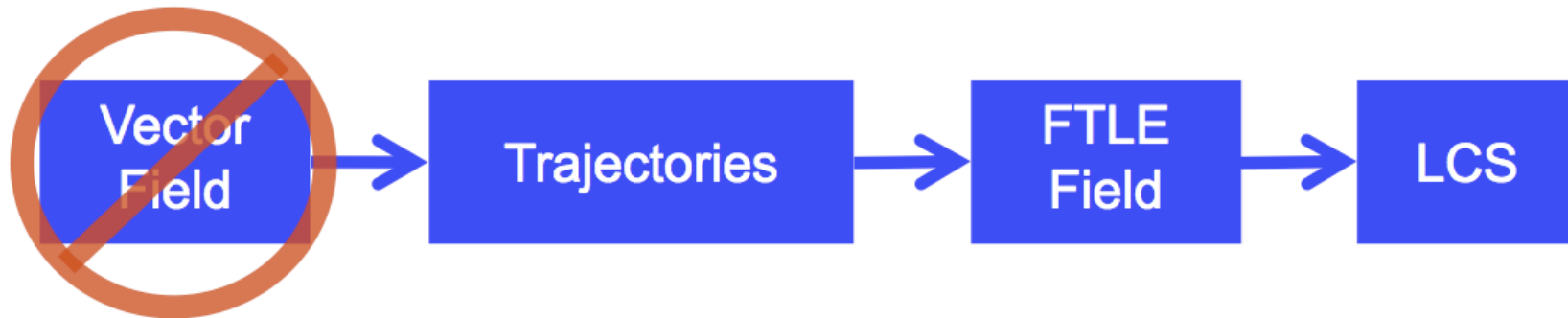
- From 4D model of bipedal walking (Norris, Marsh, Granata, Ross 2008)
- Separatrix bounds region of recovery around walking solution

- Outside of this “tube”, falling occurs



Separatrices: biological phenomena

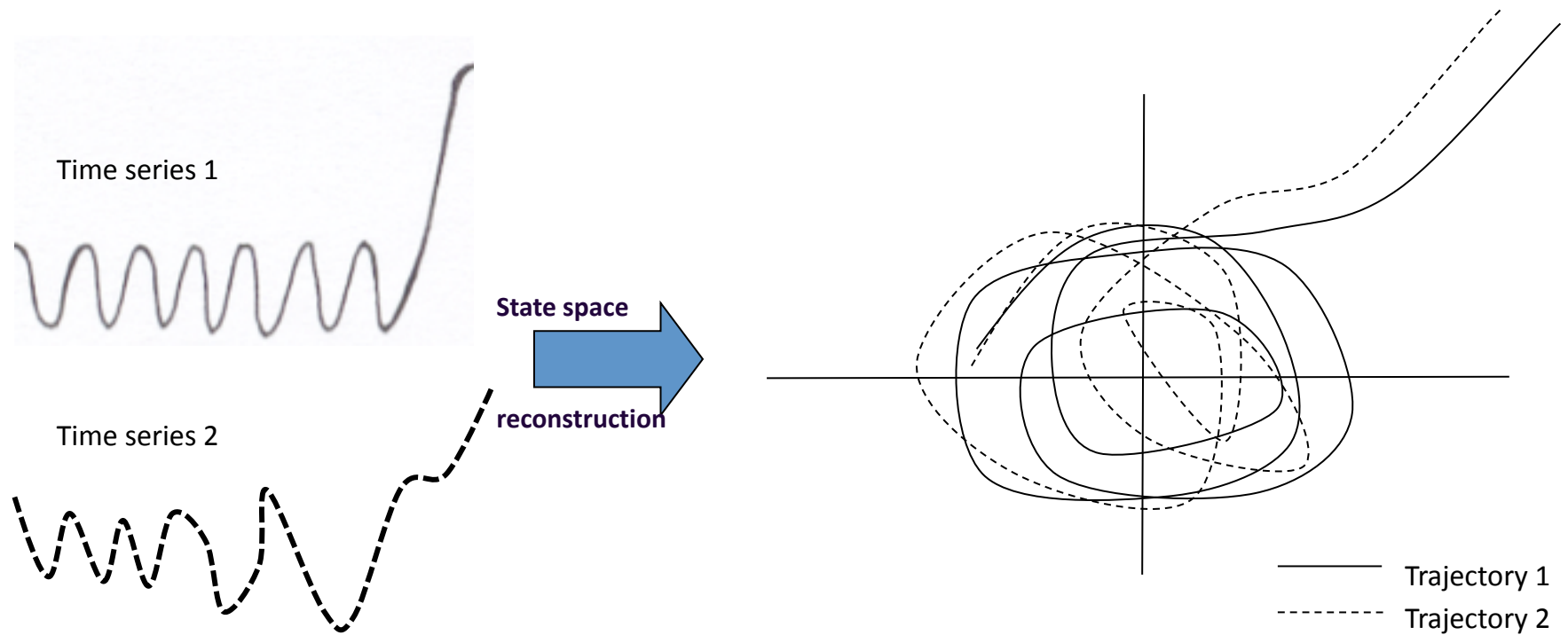
- Structure *even when no vector field known*
- For example, experimental time series data



Refs: Tanaka & Ross (2008), *Nonlinear Dynamics*

Tanaka, Nussbaum, Ross (2008), *Journal of Biomechanics*

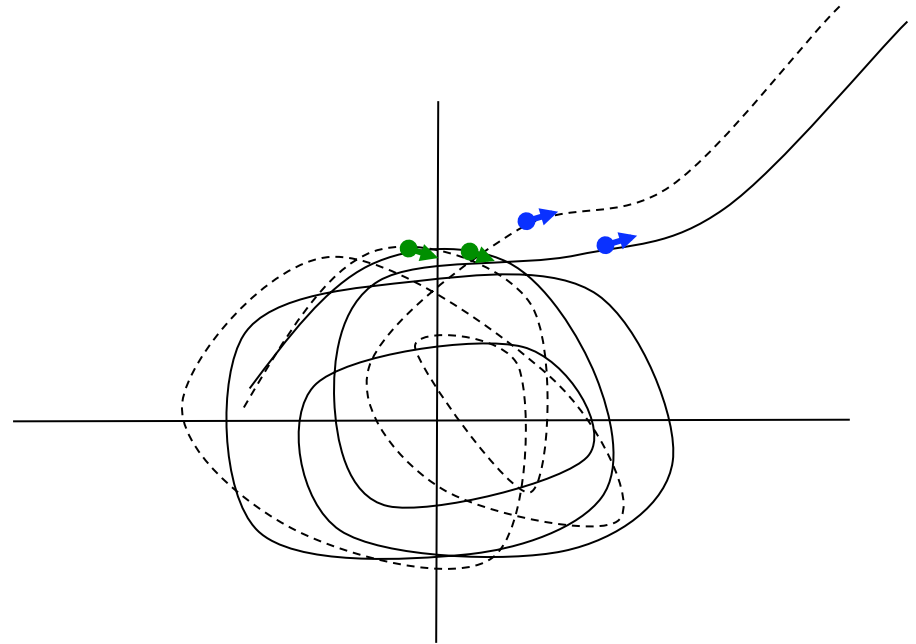
Time series to state space structure



- From time-series data to state space trajectories

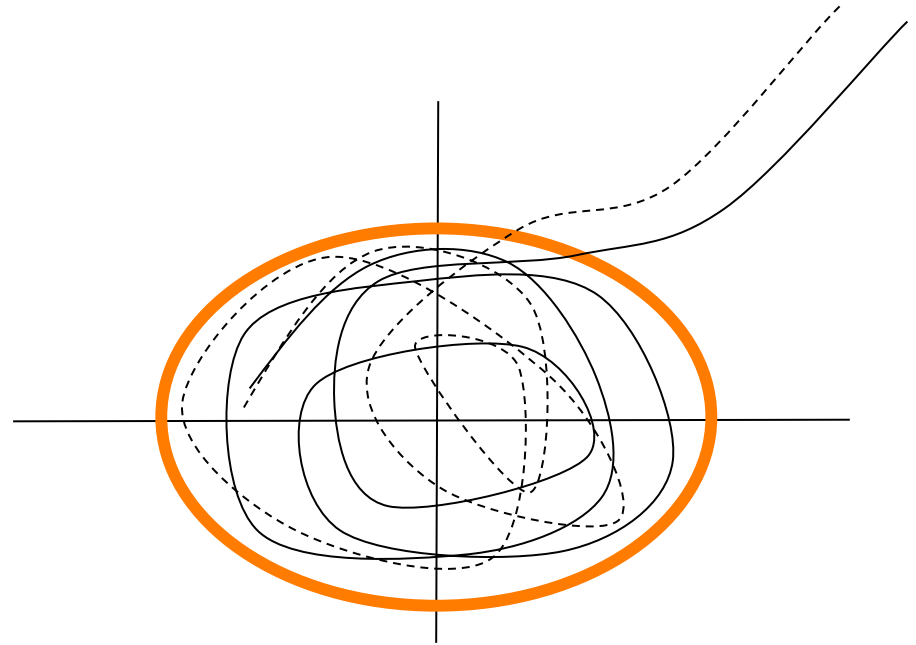
Time series to state space structure

- **Two green points have small divergence**
- **Two blue points have small divergence**
- **Green and blue points diverge rapidly**



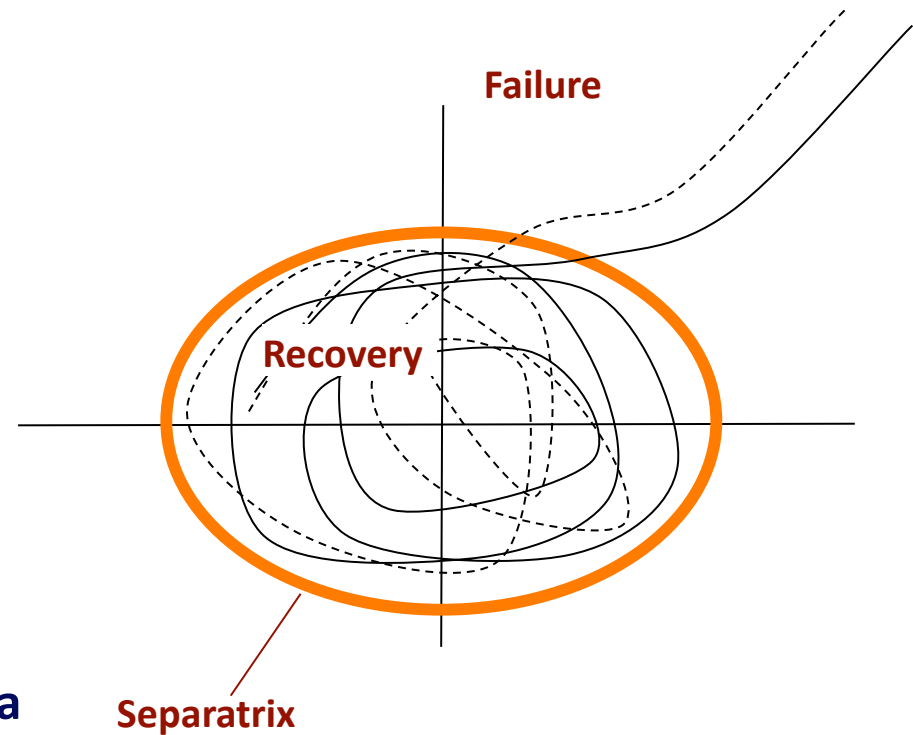
Time series to state space structure

- **Two green points have small divergence**
- **Two blue points have small divergence**
- **Green and blue points diverge rapidly**
- **A thin region of high divergence**



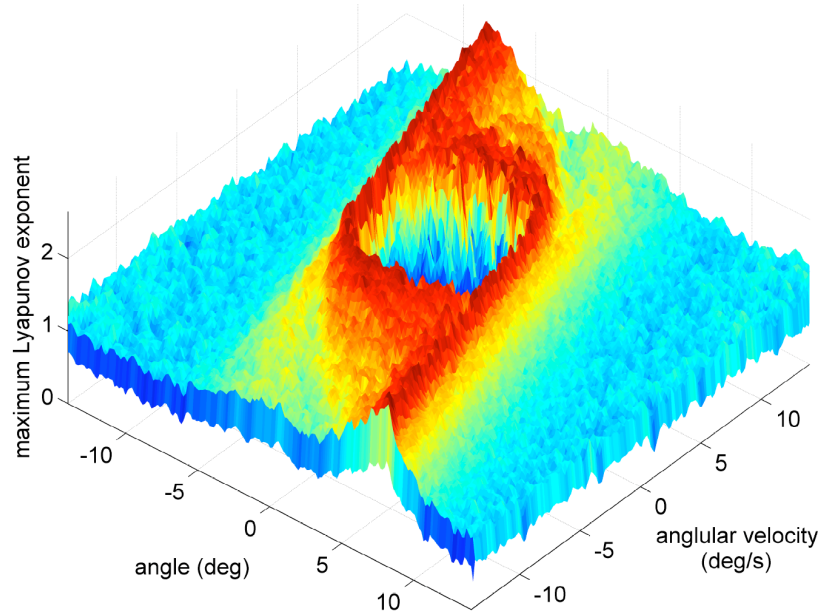
Time series to state space structure

- Two green points have small divergence
- Two blue points have small divergence
- Green and blue points diverge rapidly
- A thin region of high divergence
- This dynamical boundary separates a region of recovery from one of failure



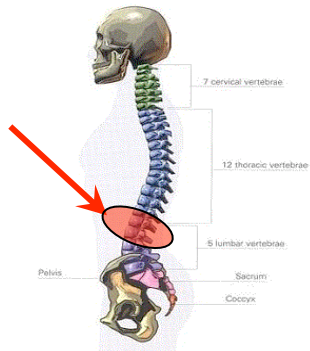
Comparison with other methods

- Other methods obtaining divergence measures from time series:
 - Yield only one number (*a state-space average*)
 - Assume an attracting set exists *and there may not be one*
 - E.g., Wolf et al 1985, Sano & Sawada 1985, Kantz & Schreiber 2004
- This method yields a sensitivity field, yielding important state space structure
 - Even for noisy, messy experimental data
 - Or mechanical systems with no attractors



Torso instability is often associated with LBP

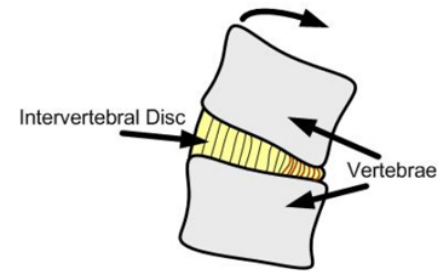
Local Instability



Torso Instability

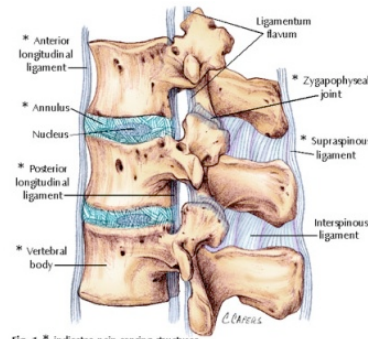
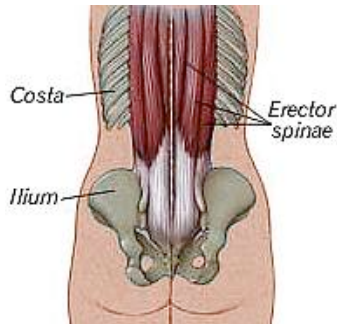
Excessive Strain

>10° Rotation



Injury

Low Back Pain



Basins of recovery in balance control

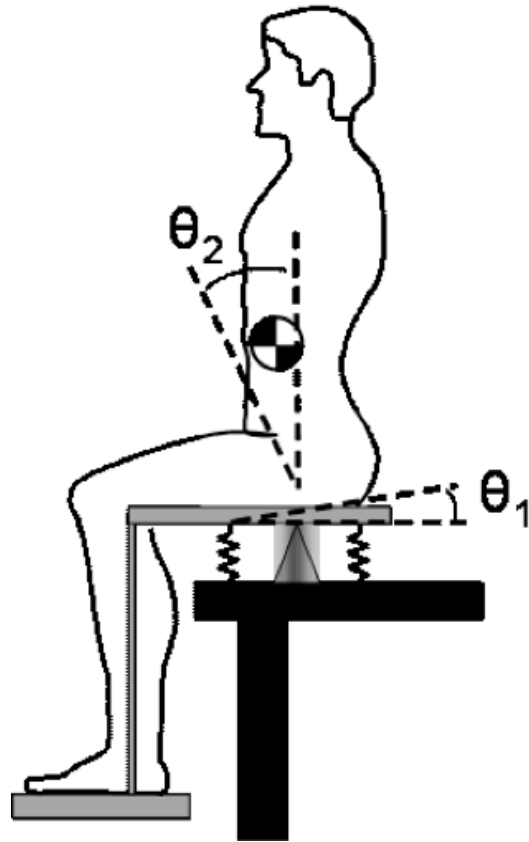
- **Wobble chair (measure torso stability, linked to low back pain)**



- **Movement of lumbar spine maintains stability**

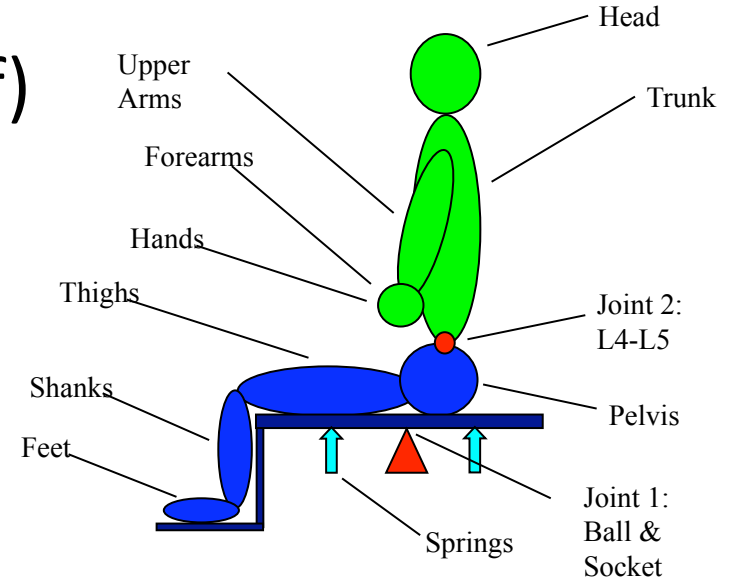
Wobble chair

- Wobble chair schematic (slice through fore-aft plane)



Wobble Chair Model (2 dof)

- Anthropometric data was used to calibrate the model
- System reduced to a double inverted pendulum (2D CS 4D SS)
- Solved using Lagrange's equations



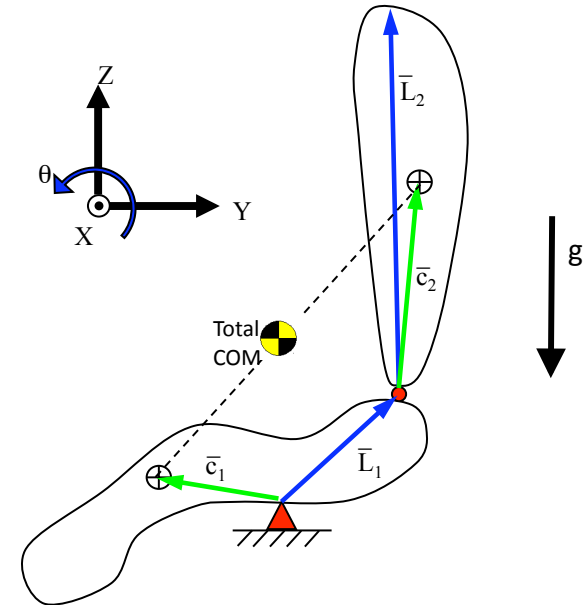
$$M\ddot{\theta} + C\dot{\theta}^2 + G(\theta) = \tau$$

$$M = \begin{bmatrix} m_1 \|\bar{c}_1\|^2 + m_2 \|\bar{L}_1\|^2 + I_1 & m_2 (R'_{\theta_1} \bar{L}_1) \cdot (R'_{\theta_2} \bar{c}_2) \\ m_2 (R'_{\theta_1} \bar{L}_1) \cdot (R'_{\theta_2} \bar{c}_2) & m_2 \|\bar{c}_2\|^2 + I_2 \end{bmatrix}$$

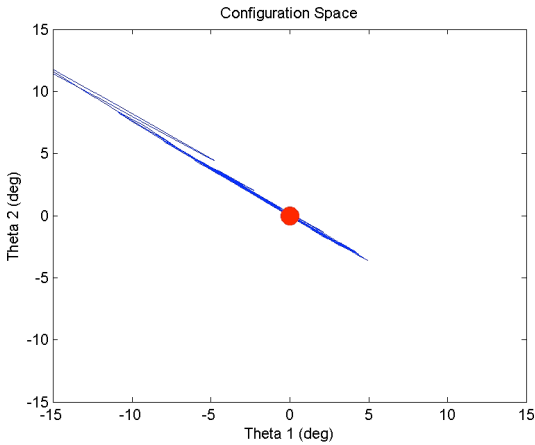
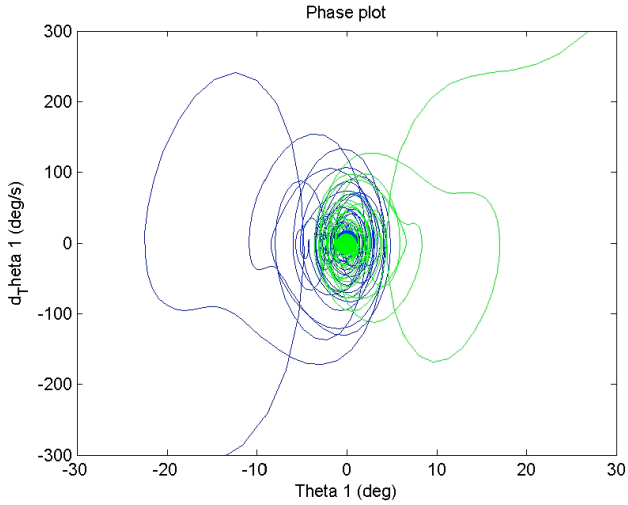
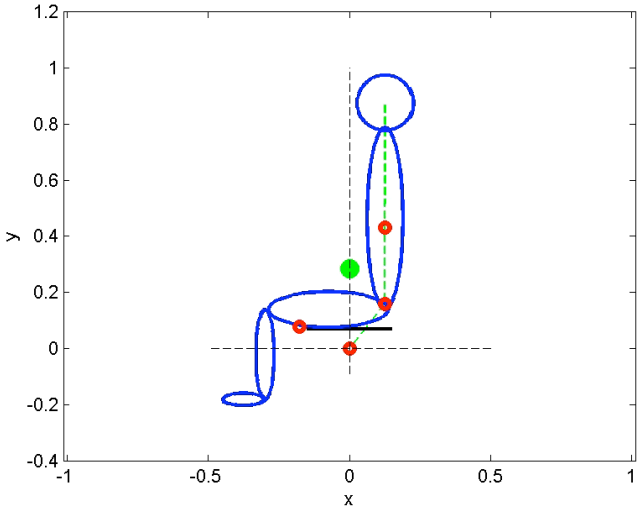
$$C = \begin{bmatrix} 0 & m_2 (R'_{\theta_1} \bar{L}_1) \cdot (R''_{\theta_2} \bar{c}_2) \\ m_2 (R''_{\theta_1} \bar{L}_1) \cdot (R'_{\theta_2} \bar{c}_2) & 0 \end{bmatrix}$$

$$G = \begin{bmatrix} m_1 \bar{g} \cdot (R'_{\theta_1} \bar{c}_1) + m_2 \bar{g} \cdot (R'_{\theta_1} \bar{L}_1) \\ m_2 \bar{g} \cdot (R'_{\theta_2} \bar{c}_2) \end{bmatrix}$$

$$\tau = \begin{bmatrix} T_{Spr} - (T_{sk} + T_{sd} + C_p + C_d + Noise) \\ T_{sk} + T_{sd} + C_p + C_d + Noise \end{bmatrix} \quad \theta = \begin{bmatrix} \theta_1 \\ \theta_2 \end{bmatrix}$$



Wobble Chair Simulation



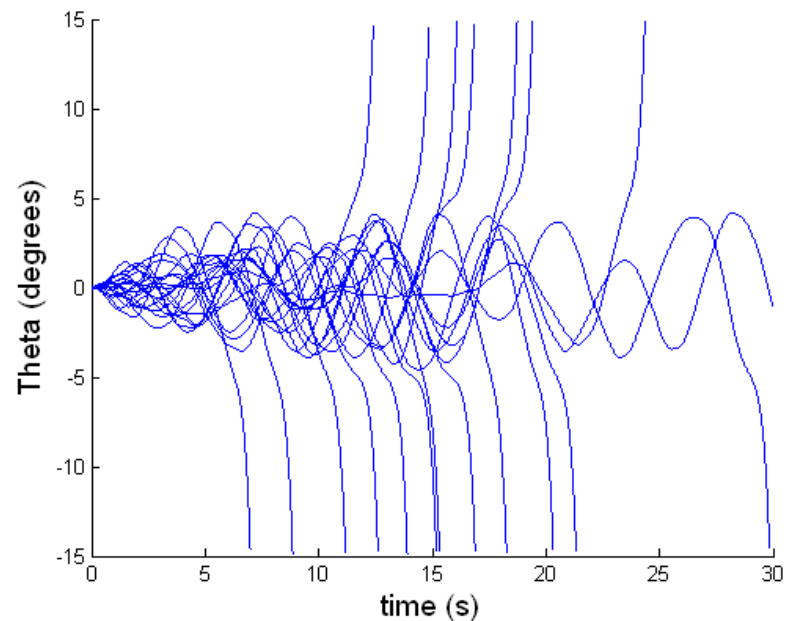
Experimental trials: sampling state space



Previous wobble chair experiments

[Tanaka and Granata 2007, Lee and Granata 2008]

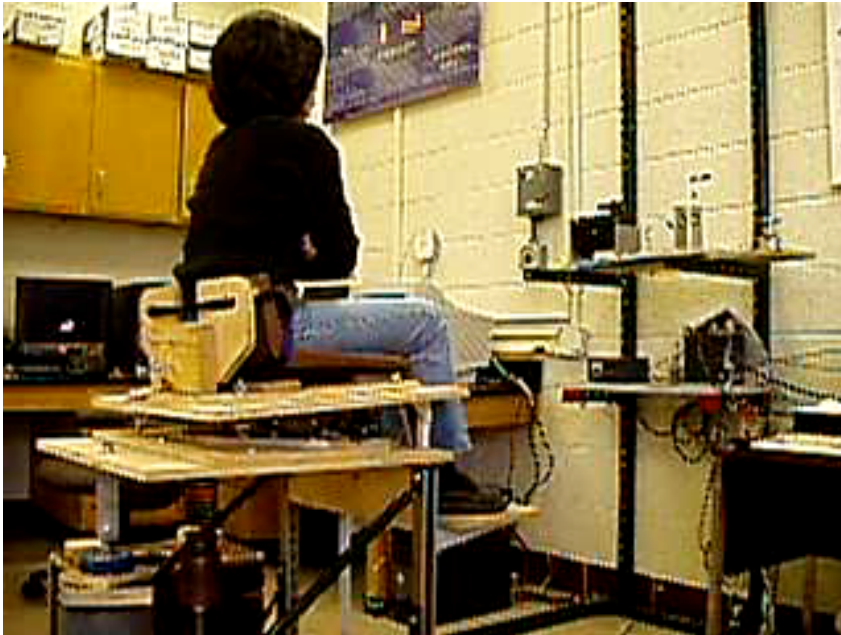
- Averaged over the entire time series
- Result was a single scalar value λ_{\max}



Measure orientation time history for multiple trials

Attempt to identify boundaries

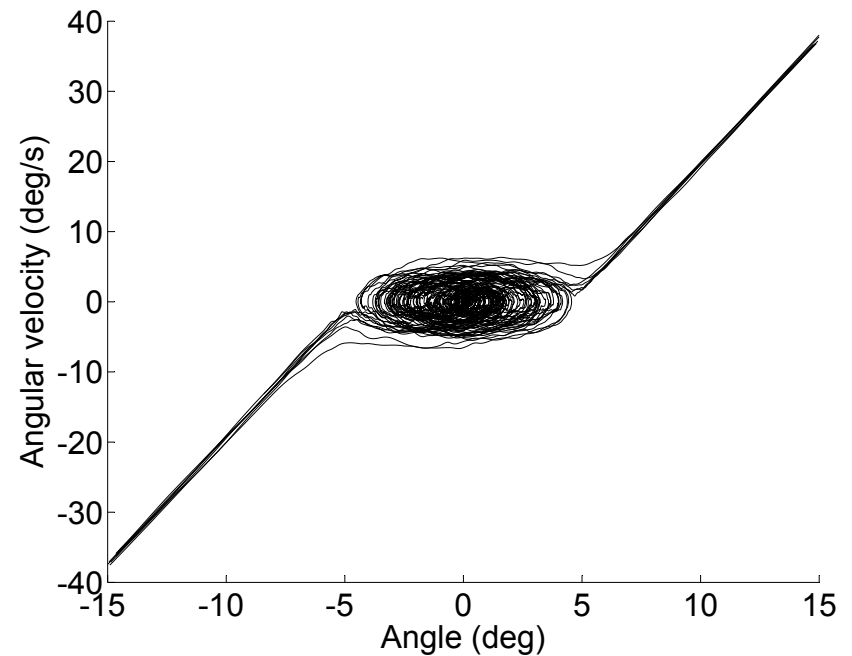
Experimental trials: sampling state space



Previous wobble chair experiments

[Tanaka and Granata 2007, Lee and Granata 2008]

- Averaged over the entire time series
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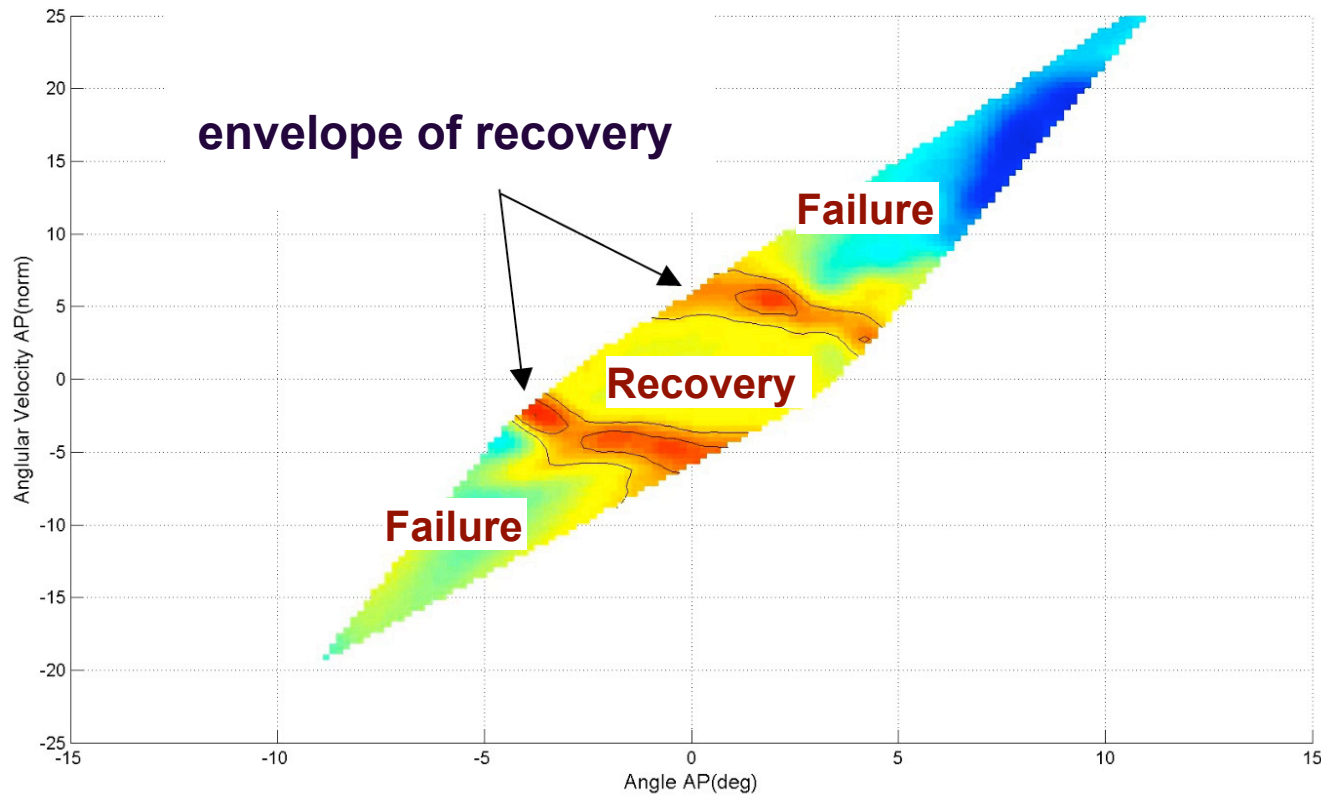


Measure orientation time history for multiple trials

Attempt to identify boundaries

Boundary of basin of recovery

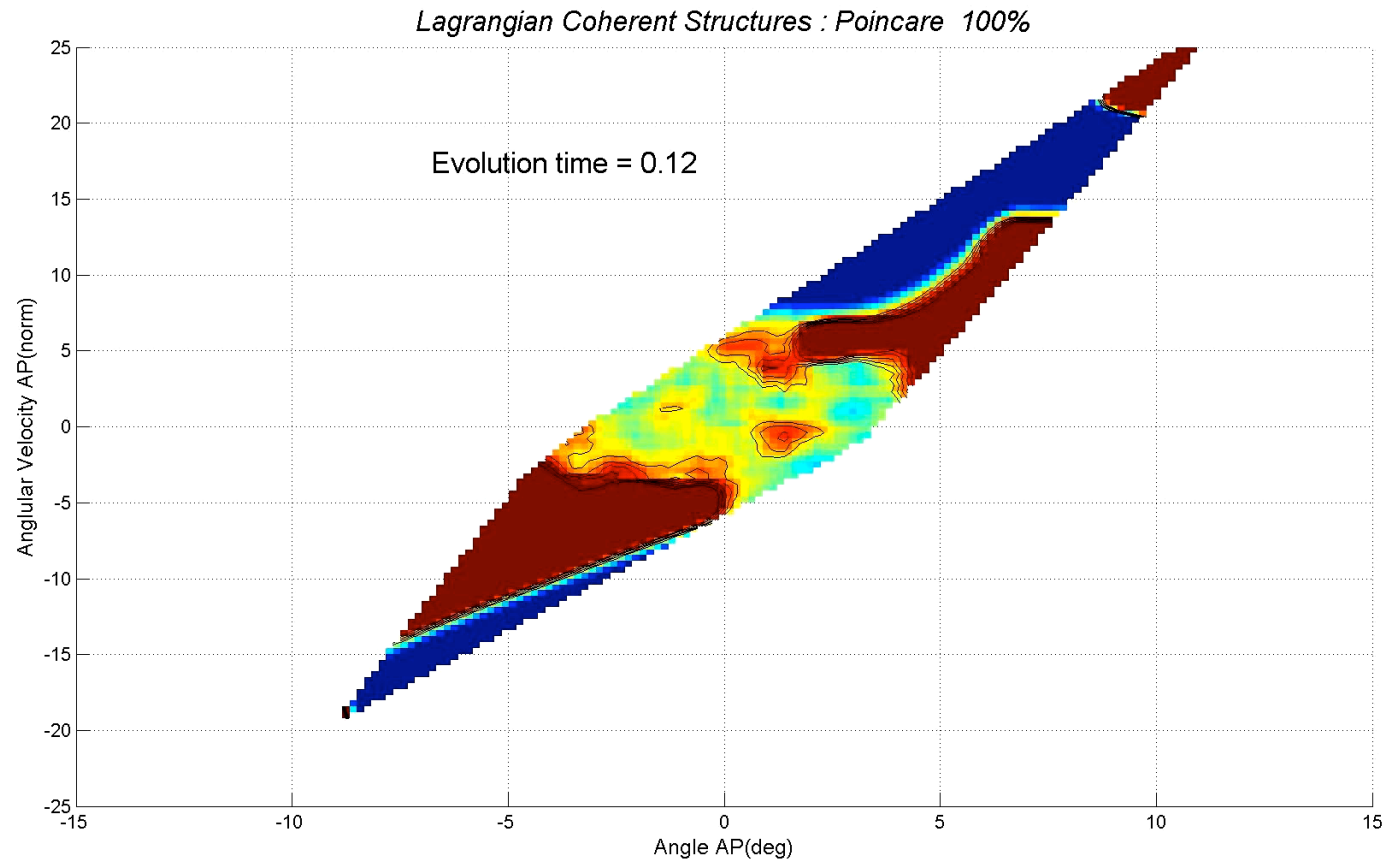
- Measure motion
- Time-series data to state space trajectories, even sparse data
- Compute the recovery/failure boundary, the envelope of recovery



Basin of recovery bounded by ridges in sensitivity field

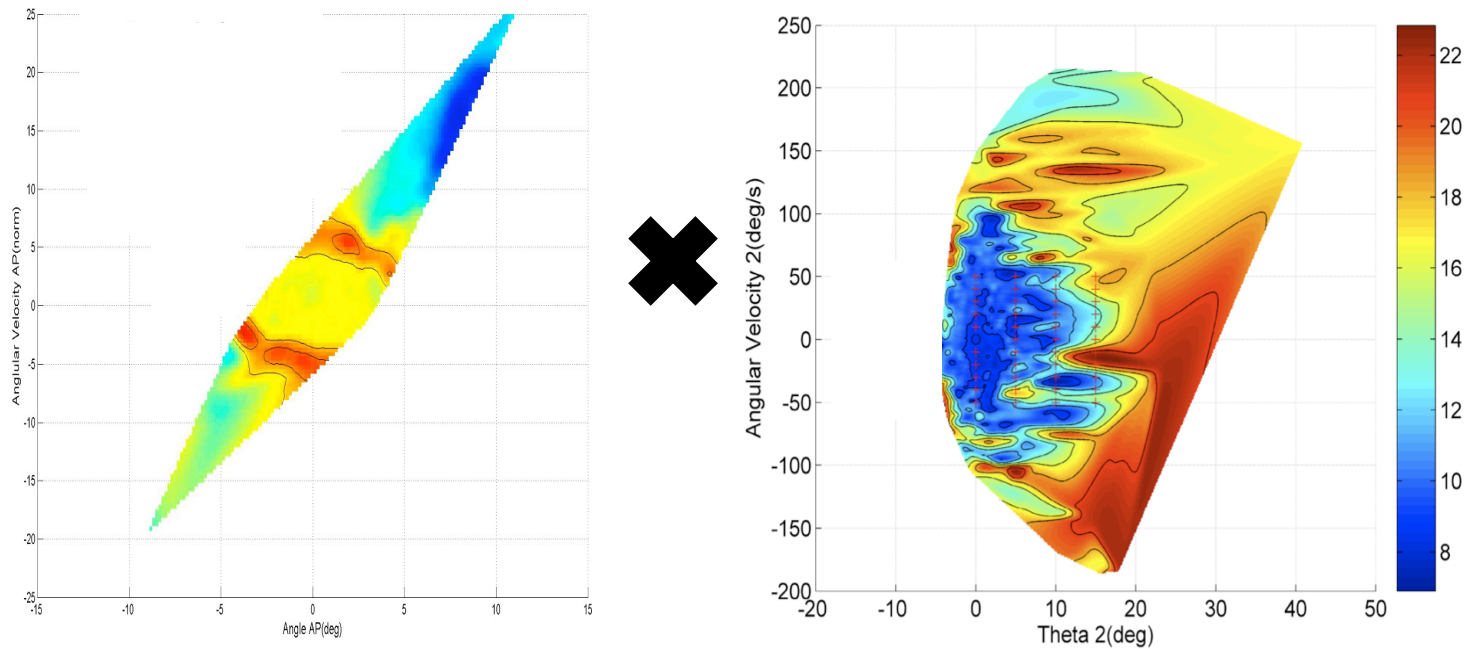
- **Sensitivity field** could be max finite-time Lyapunov exponent field
- Measures **divergence** between trajectories starting near a point
- Ridges are **boundaries** between qualitatively different trajectories; originally developed as “Lagrangian coherent structures” in fluid mechanics

Sensitivity field reveals structure as **evolution time** increases



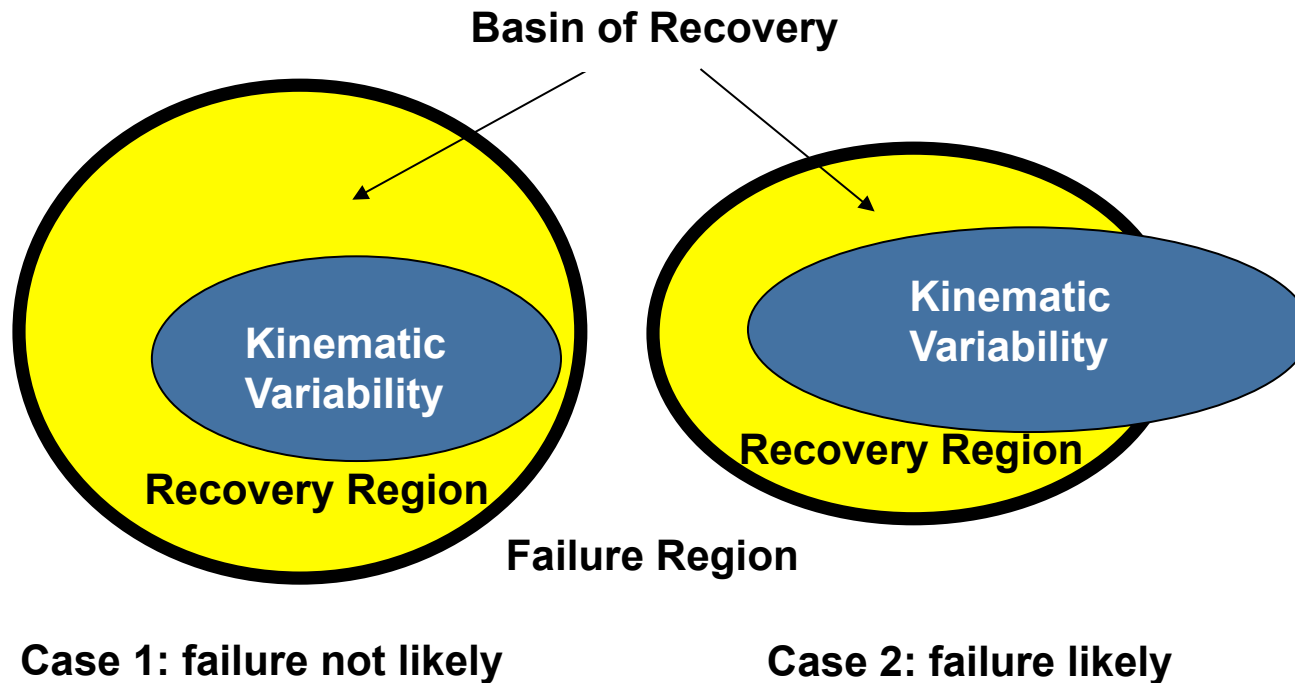
Boundaries in higher dimension

- Basin of recovery in n-dim system bounded by (n-1)-dim surface
- 2 angles and their rates (2D X 2D = 4D system)
- Hard to visualize, but can compute *4D basin bounded by 3D surface*
- Basin size affected by environmental conditions



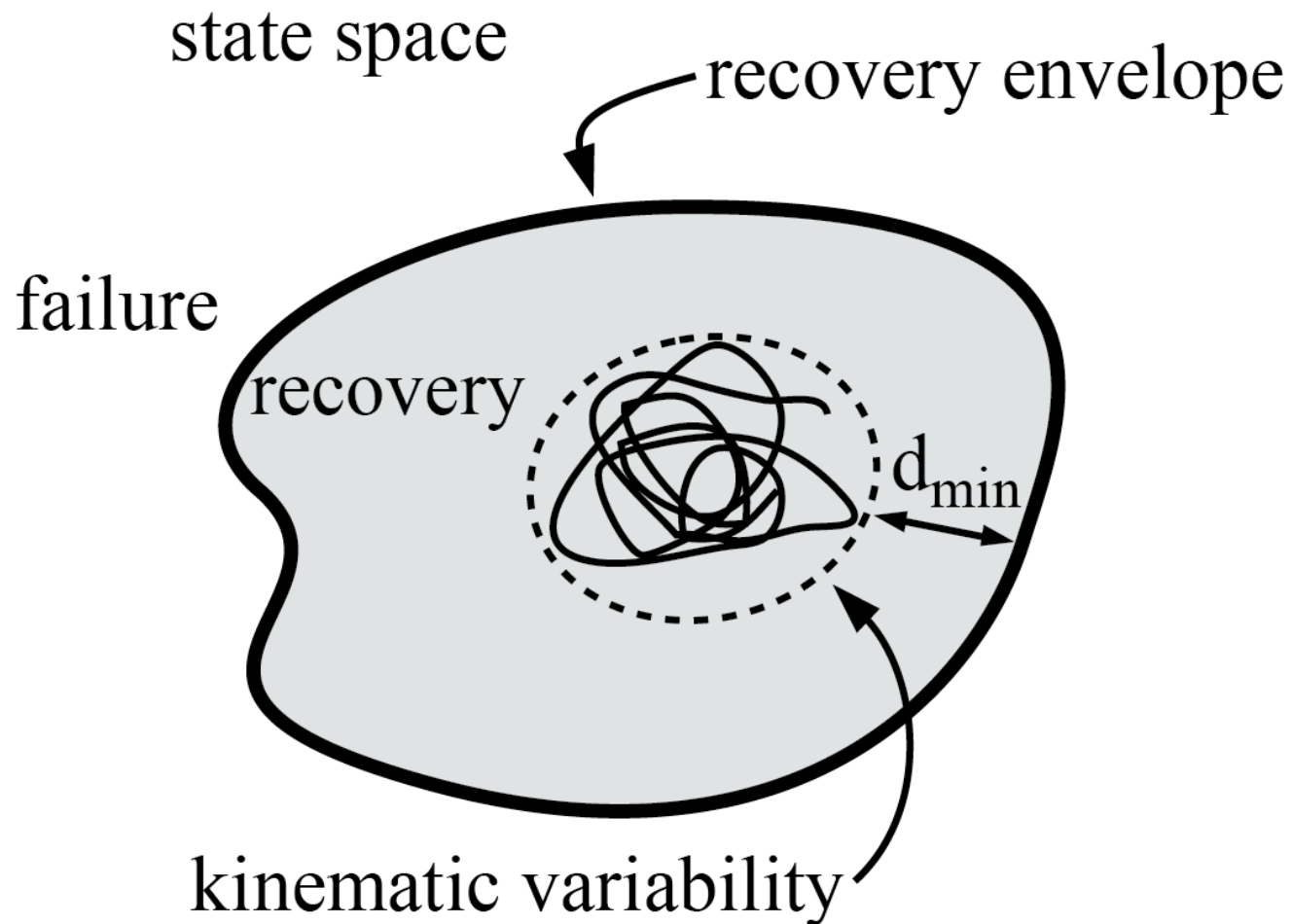
Recovery envelope: the basin boundary

- Size and shape of basin is function of *possible* natural dynamics
- Kinematic variability: currently explored region of state space
- When kinematic variability exceeds basin of recovery => failure



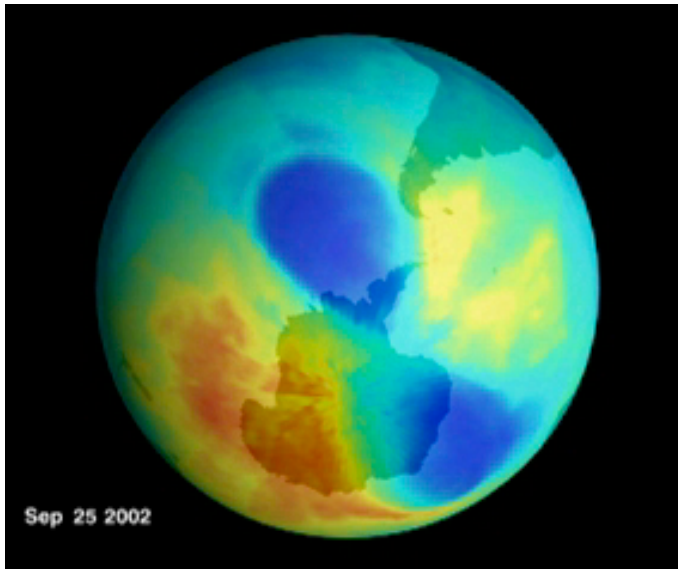
A new tool in the evaluation of risk of failure?

- **Minimum distance between kinematic variability and recovery envelope can measure risk of failure**

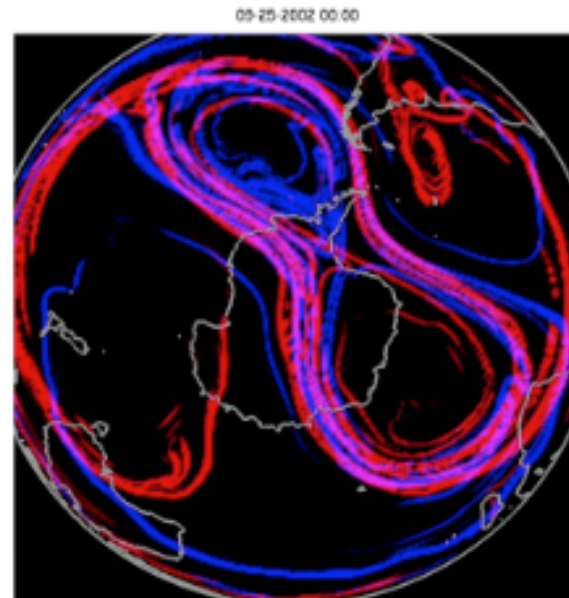


Separatrices (LCS)

- Structure of trajectories and transport in chaotic, time-dependent vector field



Ozone concentration



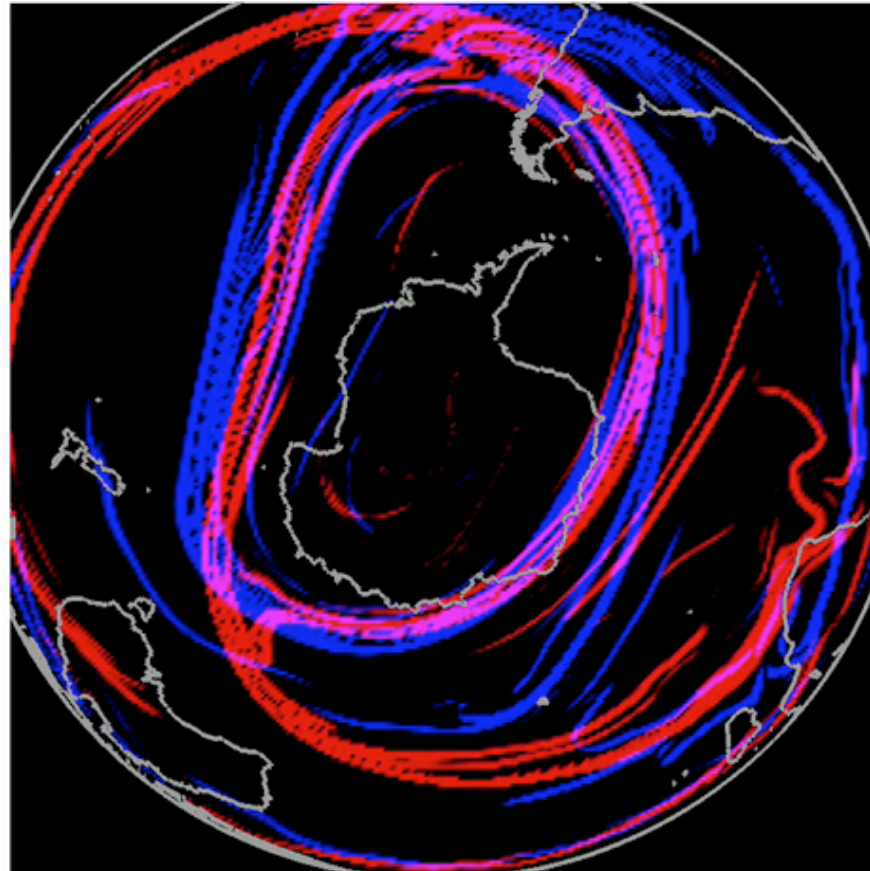
Separatrices from wind data

2002 Ozone Hole Splitting Event

Atmospheric transport barriers

Are there isolated “aero-ecosystems”?

09-09-2002 06:00



Atmospheric transport of pathogens

Determine atmospheric pathways

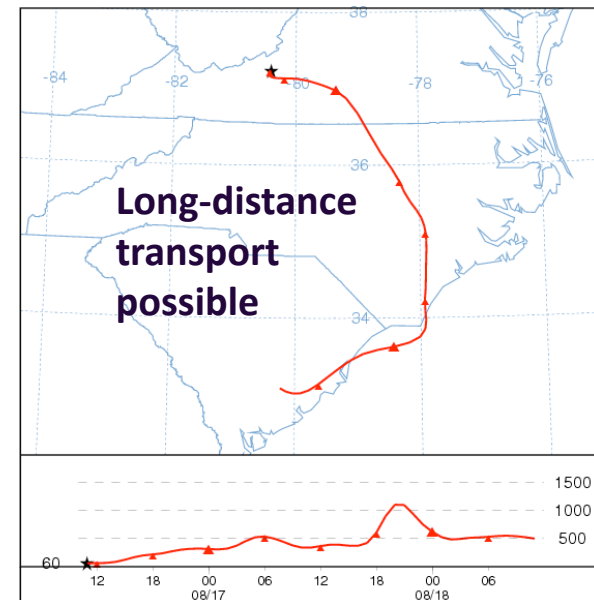
- For airborne biological pathogens which lead to spread of plant disease

Relevant scales

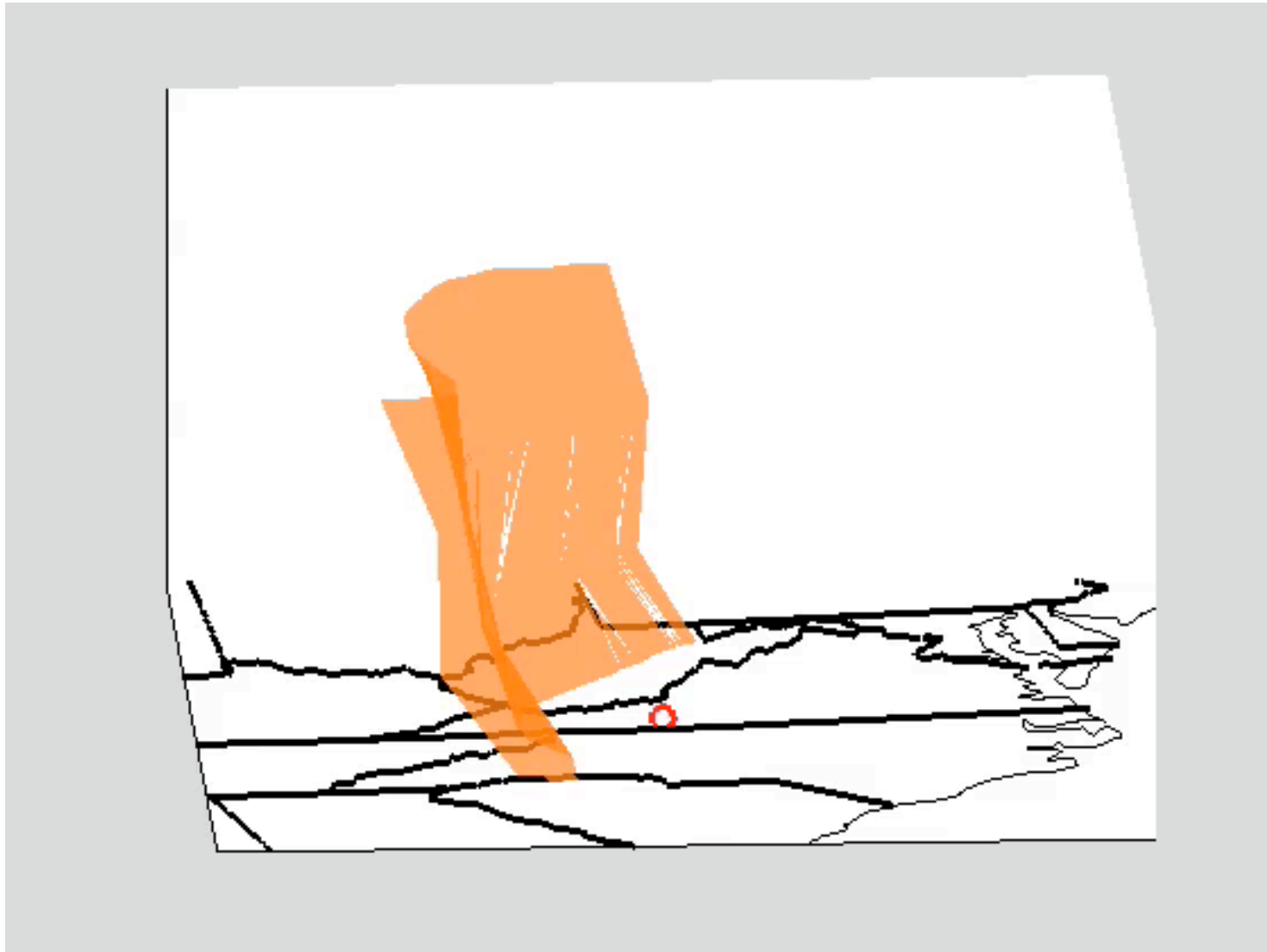
- Horizontal spread over large areas
- Pathogens may separate from fluid flow and change shape or clump



Crop diseases spread by airborne spores



Atmospheric transport barriers (3D)



Outlook

- **Dynamical boundaries/separatrices reveal state space structure displayed by trajectories; even in noisy, experimental time series data**
- **Ridges in finite-time Lyapunov exponent field over sampled state space**
 - **Field average gives usual Lyapunov exponent quoted by others**
- **Applications: torso stability, walking stability, bipedal robotics, prosthetics, ship capsizing**
- **Limitations**
 - **Applicable only to mechanical systems?**
 - **Data requirements? Quantity and quality of data to see structure?**
 - **High dimensions difficult; choose appropriate coarse variables**

